

# Lecture 2: Variational Problems II

Note Title

## ★ The Euler-Lagrange Eqn.

Consider a more general problem than the particle system in Lecture 1.

Let  $\Gamma := \{y \in C^1[x_1, x_2] \mid y(x_i) = y_i, i=1,2\}$

Find  $y \in \Gamma$  s.t.

$$(*) \quad I = \int_{x_1}^{x_2} f(x, y, y') dx \rightarrow \min!$$

where  $f \in C^2$  in each variable.

As before, assume  $E: y = y^*(x)$  is the extremal (minimizer) and

consider a one-parameter variation of  $E$ , i.e.,  $E_\varepsilon: y = y^*(x, \varepsilon) = y_\varepsilon^*(x)$

with  $\frac{\partial y^*}{\partial x}, \frac{\partial y^*}{\partial \varepsilon}, \frac{\partial^2 y}{\partial x \partial \varepsilon}$  are all cont. fns in  $x$ ,  
and  $y^*(x, 0) = y^*(x), y^*(x_i, \varepsilon) = y_i, i=1,2$ .



given const's.



For the notational convenience, let

$$\frac{\partial y^*}{\partial \varepsilon}(x, \varepsilon) =: \eta(x, \varepsilon). \text{ Then clearly}$$

$$\eta(x_i, \varepsilon) = 0, \quad i=1,2.$$

Putting these in (\*), we get

$$I(\varepsilon) := \int_{x_1}^{x_2} f(x, y^*(x, \varepsilon), y^{*\prime}(x, \varepsilon)) dx$$

For (\*) to be min.,  $\delta I = \frac{dI}{d\varepsilon}(0) = 0$  is a must.

$$\frac{dI}{d\varepsilon} = \int_{x_1}^{x_2} \left( \frac{\partial f}{\partial y}(x, y_\varepsilon^*, y_\varepsilon^{*\prime}) \cdot \frac{\partial y_\varepsilon^*}{\partial \varepsilon} + \frac{\partial f}{\partial y'}(x, y_\varepsilon^*, y_\varepsilon^{*\prime}) \cdot \frac{\partial y_\varepsilon^{*\prime}}{\partial \varepsilon} \right) dx$$

$$= \int_{x_1}^{x_2} \left( \frac{\partial f}{\partial y}(x, y_\varepsilon^*, y_\varepsilon^{*\prime}) \cdot \eta + \frac{\partial f}{\partial y'}(x, y_\varepsilon^*, y_\varepsilon^{*\prime}) \cdot \eta' \right) dx$$

Int.  
by  
parts

$$= \int_{x_1}^{x_2} \frac{\partial f}{\partial y}(x, y_\varepsilon^*, y_\varepsilon^{*\prime}) \eta dx$$

$$+ \underbrace{\eta \frac{\partial f}{\partial y'}(x, y_\varepsilon^*, y_\varepsilon^{*\prime}) \Big|_{x_1}^{x_2}} - \int_{x_1}^{x_2} \eta \frac{d}{dx} \frac{\partial f}{\partial y'}(x, y_\varepsilon^*, y_\varepsilon^{*\prime}) dx$$

$$= \int_{x_1}^{x_2} \eta \left\{ \frac{\partial f}{\partial y}(x, y_\varepsilon^*, y_\varepsilon^{*\prime}) - \frac{d}{dx} \frac{\partial f}{\partial y'}(x, y_\varepsilon^*, y_\varepsilon^{*\prime}) \right\} dx$$

$$\xrightarrow{\substack{\text{set} \\ \varepsilon=0}} \int_{x_1}^{x_2} \eta \left\{ \frac{\partial f}{\partial y}(x, y^*, y^{*\prime}) - \frac{d}{dx} \frac{\partial f}{\partial y'}(x, y^*, y^{*\prime}) \right\} dx$$

Since  $\eta$  is arbitrary as long as  $\eta \in \Gamma_0$ ,  
by the F.L.C.V., we must have admissible variation  
 $P$  &  $\eta(x_i) = 0$

$$\frac{\partial f}{\partial y} - \frac{d}{dx} \frac{\partial f}{\partial y'} = 0$$

The Euler-Lagrange Eqn.

Note that we assumed  $\eta' \in C[x_1, x_2]$  for Int. by Parts & exchange of  $\frac{\partial}{\partial \varepsilon}, \frac{\partial}{\partial x}$  to be valid.

Also,  $\frac{d}{dx} \frac{\partial f}{\partial y'}$  must be in  $C[x_1, x_2]$ .

$$\frac{d}{dx} \frac{\partial f}{\partial y'} = \frac{\partial^2 f}{\partial y' \partial x} + \frac{\partial^2 f}{\partial y' \partial y} y' + \frac{\partial^2 f}{\partial y'^2} y''$$

$$\Rightarrow y'' \in C[x_1, x_2], \text{ i.e., } y \in C^2[x_1, x_2].$$

Also, we will assume  $\frac{\partial^2 f}{\partial y'^2} \neq 0$  on  $[x_1, x_2]$ .

If  $\frac{\partial^2 f}{\partial y'^2} \equiv 0$  on  $[x_1, x_2]$ , a regular variational problem

then it's called

a singular variational problem.

In this case, the derived differential eqn. will be of the first order, not the second.

So far, in deriving the E-L eqn.

the important things were:

- ① Integration by Parts + Bdry. cond.
- ② F. L. C. V.

\* Fundamental Lemma of the Calc. of Var.

If  $M \in C[x_1, x_2]$ ,  $\eta \in C^1[x_1, x_2]$  with  $\eta(x_i) = 0, i=1, 2$ , and  $\int_{x_1}^{x_2} \eta(x) M(x) dx = 0$  - (\*)

for all such  $\eta$ , then  $M(x) \equiv 0$  on  $[x_1, x_2]$ .

Again, this is a necessary cond.

(Proof) Assume  $M(x) \neq 0$  on  $x \in [x_1, x_2]$ .

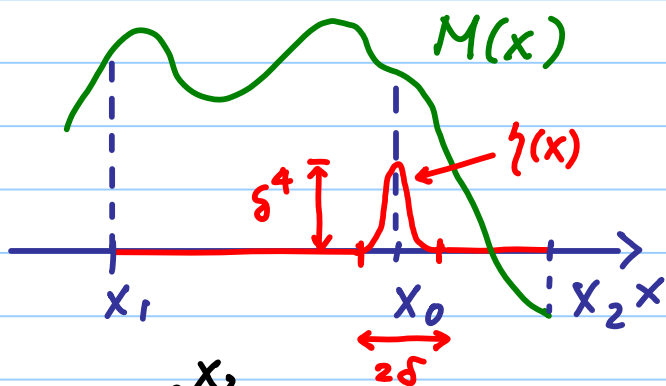
Then,  $\exists x_0 \in [x_1, x_2]$  s.t.  $M(x_0) \neq 0$ .

WLOG, suppose  $M(x_0) > 0$ .

$M \in C[x_1, x_2] \Rightarrow \exists \delta > 0$  s.t.  $M(x) > 0$   
 $\forall x$  in  $|x - x_0| < \delta$ .

Now consider the following fcn  $\eta$ :

$$\eta(x) := \begin{cases} 0 & \text{if } |x - x_0| > \delta \\ (x - x_0 + \delta)^2 (x - x_0 - \delta)^2 & \text{if } |x - x_0| \leq \delta. \end{cases}$$



This  $\eta$  satisfies the cond., i.e.,  $\eta \in C^1[x_1, x_2]$  with  $\eta(x_i) = 0, i=1, 2$ . But  $\eta \notin C^2[x_1, x_2]$ . Also  $\eta(x) > 0$  for  $|x - x_0| < \delta$ .

$$\text{Now, } \int_{x_1}^{x_2} \eta(x) M(x) dx = \int_{x_0-\delta}^{x_0+\delta} (x-x_0+\delta)^2 (x-x_0-\delta)^2 M(x) dx > 0$$

This contradicts the premise (\*). #

\* A generalization to  $n$  unknown funcs

Suppose  $y_k = y_k(x), k=1, \dots, n$

$y_k \in C^1[x_1, x_2], y_k(x_i) = y_k^{(i)}, k=1:n, i=1, 2.$

$$I = \int_{x_1}^{x_2} f(x, y_1, \dots, y_n, y_1', \dots, y_n') dx$$

where  $f$  is in  $C^2$  in each variable.

Then using the same argument, we can derive  $n$  Euler-Lagrangé eqn's.

$$\frac{\partial f}{\partial y_k} - \frac{d}{dx} \frac{\partial f}{\partial y_k'} = 0, k=1, \dots, n.$$

Ex. Newton's eqn's of Motion in  $\mathbb{R}^3$ , i.e.,  $n=3$ .

$$f = \frac{m}{2} (\dot{x}^2 + \dot{y}^2 + \dot{z}^2) - V(x, y, z)$$

Here,  $(x, y, z, t) \leftrightarrow (y_1, y_2, y_3, x)$   
 $\cdot \leftrightarrow \frac{d}{dx}$

$$\text{So, } \frac{\partial f}{\partial x} - \frac{d}{dt} \frac{\partial f}{\partial \dot{x}} = 0 \iff -\frac{\partial V}{\partial x} - m\ddot{x} = 0$$

$$\iff m\ddot{x} = -\frac{\partial V}{\partial x}$$

Similarly for  $y$  &  $z$ ,  
 we get  $m \dot{r}(t) = -\nabla V(r(t))$

Now, how about the same system  
 in  $\mathbb{R}^3$ , but in the polar coordinates  
 $(r, \theta, \varphi)$ ?  $\Rightarrow$  HW problem!

$$\left\{ \begin{array}{l} \frac{\partial f}{\partial r} - \frac{d}{dt} \frac{\partial f}{\partial \dot{r}} = 0 \\ \frac{\partial f}{\partial \theta} - \frac{d}{dt} \frac{\partial f}{\partial \dot{\theta}} = 0 \\ \frac{\partial f}{\partial \varphi} - \frac{d}{dt} \frac{\partial f}{\partial \dot{\varphi}} = 0 \end{array} \right. \Rightarrow ??$$