UCRL-JC-122065 PREPRINT

A Numerical Study of Shock-Acceleration of a Diffuse Helium Cylinder

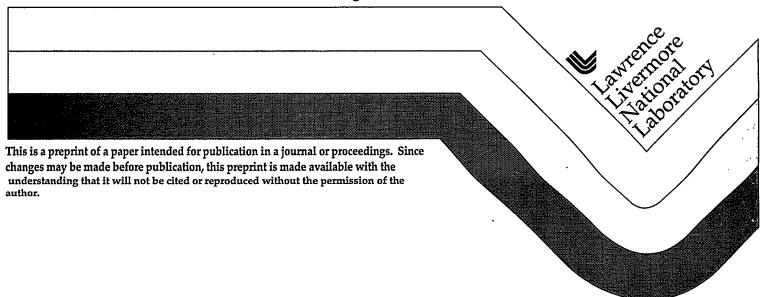
Jeffrey A. Greenough, John Bell, Phillip Colella Center for Computational Sciences & Engineering 7 1 3 1995 P.O. Box 808, L-316 Lawrence Livermore National Laboratory Livermore, CA 94551 U.S.A.

Elbridge G. Puckett
Department of Mathematics and Institute of Theoretical Dynamics
University of California
Davis, CA 95616 U.S.A.

Jeffrey W. Jacobs
Department of Aerospace and Mechanical Engineering
University of Arizona
Tucson, AZ 85721 U.S.A.

This was prepared for submittal to the Proceedings of the 20th International Symposium on Shock Waves (ISSW), Pasadena, California July 23-28, 1995

August 1995



DISCLAIMER

This document was prepared as an account of work sponsored by an agency of the United States Government. Neither the United States Government nor the University of California nor any of their employees, makes any warranty, express or implied, or assumes any legal liability or responsibility for the accuracy, completeness, or usefulness of any information, apparatus, product, or process disclosed, or represents that its use would not infringe privately owned rights. Reference herein to any specific commercial product, process, or service by trade name, trademark, manufacturer, or otherwise, does not necessarily constitute or imply its endorsement, recommendation, or favoring by the United States Government or the University of California. The views and opinions of authors expressed herein do not necessarily state or reflect those of the United States Government or the University of California, and shall not be used for advertising or product endorsement purposes.

DISCLAIMER

Portions of this document may be illegible in electronic image products. Images are produced from the best available original document.

A Numerical Study of Shock-Acceleration of a Diffuse Helium Cylinder *

Jeffrey A. Greenough, John B. Bell, Phillip Colella Center for Computational Sciences and Engineering Lawrence Livermore National Laboratory Livermore, CA 94550

Elbridge G. Puckett Department of Mathematics and Institute of Theoretical Dynamics University of California Davis, CA 95616

Jeffrey W. Jacobs Department of Aerospace and Mechanical Engineering University of Arizona Tucson, AZ 85721

Abstract: The development of a shock-accelerated diffuse Helium cylindrical inhomogeneity is investigated using a new numerical method. The new algorithm is a higher-order Godunov implementation of the so-called multi-fluid equations. This system correctly models multiple component mixtures by accounting for differential compressibility effects. This base integrator is embedded in an implementation of adaptive mesh refinement (AMR) that allows efficient increase in resolution where the computational effort is concentrated where high accuracy, or increased resolution, are required. Qualitative and quantitative comparison with previous experimental data is excellent. The simulations show that counter-sign vortex blobs are deposited in the jet core by baroclinic generation of the curved shock wave as it traverses the jet. This vorticity deposition occurs over timescales that scale with the shock passage time ($\sim 10\mu \rm sec$). Three phases of development are identified and characterized. The first is the weak deformation (WD) phase, where there is weak distortion of the Helium jet due to weak vorticity induced velocity effects. The second phase is the strong deformation (SD) phase where there is large distortion for the jet and the vortex blobs due to large induced velocity effects. The last is a relaxation/reorganization (RR) phase where the vorticty field reorganizes into

^{*}Work performed under the auspices of the U.S. Department of Energy by the Lawrence Livermore National Laboratory under contract W-7405-Eng-48. Additional support under contract W-7405-Eng-48 by the Applied Mathematical Sciences Program and the HPCC Grand Challenge Program of the Office of Scientific Computing at DOE, Defense Nuclear Agency under IACRO 94-831 and the Center for Advanced Fluid Dynamics Applications under contract B295120.



point-like vortex pair.

Key words: Shock-Accelerated diffuse interfaces, multi-fluid methods, adaptive methods

1. Introduction

The evolution of interfaces accelerated by shock waves has been the focus of investigations for many years. This class of problem has applications in such disparate fields as inertial confinement fusion (ICF) and high-speed combustion. The literature is rich in studies on the stability of plane interfaces [1], [2], [3], and [4]. The stability of curved interfaces has recieved similar attention culminating with the study of [5], where cylindrical or spherical shapes, constructed from soap bubbles or nitrocellulose membranes, where filled with either a light or heavy gas and then accelerated by a shock wave. This last study and the previous works relied on such membranes to separate gases. A different class of experiment on acceleration of gaseous cylinders was developed by [7] and [6] that was free from the effect of membranes and used a planar laser induced fluoresence (PLIF) technique allowed for excellent visualization of the flow and quantitative information regarding mixing.

The present simulations use a high resolution adaptive method designed for multi-fluid flows with initial mixed layers. In this paper, the development of the Helium mole fraction and vorticity field will be examined. Extensive comparisons with experimental data have also been completed as well as development of a new conceptual model describing this flow, but cannot be described here due to space limitations.

2. Multi-fluid System and Numerical Method

The equations used for this study are the so-called multi-fluid equations as developed by [8] and also described in [9]. This system is based on a volume-of-fluid approach, where the fraction of a cell volume occupied by a distinguished fluid is followed during the flow evolution. The basic assumptions in this approach are that there is pressure equilibrium within a cell, there is a single velocity vector for both fluids and that all changes of state are adiabatic. The equations are given as

$$\frac{\partial f^{\alpha}}{\partial t} + \nabla \cdot (\boldsymbol{u} f^{\alpha}) = f^{\alpha} \frac{\hat{\Gamma}}{\Gamma^{\alpha}} \nabla \cdot \boldsymbol{u}$$
 (2.1)

$$\frac{\partial}{\partial t} (f^{\alpha} \rho^{\alpha}) + \nabla \cdot (\boldsymbol{u} f^{\alpha} \rho^{\alpha}) = 0$$
 (2.2)

$$\frac{\partial \rho u}{\partial t} + \nabla \cdot (u u \rho) + \nabla p = 0$$
 (2.3)

$$\frac{\partial}{\partial t} (f^{\alpha} \rho^{\alpha} E^{\alpha}) + \nabla \cdot (u f^{\alpha} \rho^{\alpha} E^{\alpha}) + p f^{\alpha} \frac{\hat{\Gamma}}{\Gamma^{\alpha}} \nabla \cdot u + f^{\alpha} \frac{\rho^{\alpha}}{\rho} u \cdot \nabla p = 0$$
 (2.4)

where f^{α} , ρ^{α} , and E^{α} are the volume fraction, density, and total energy density of fluid component α . The volume fraction is defined as $f^{\alpha} = \Lambda_{\alpha}/\Lambda$ where Λ is the volume of the cell and Λ_{α} is the volume of the cell occupied by fluid α . Γ^{α} is the sound speed gamma for fluid α , and $\hat{\Gamma} = 1/\sum_{\alpha} (f^{\alpha}/\Gamma^{\alpha})$ represents fraction weighted Γ for the mixture.

The pressure that appears in the above system is defined to be a thermodynamically consistent pressure and defined as $p = \sum_i \hat{\Gamma}(f^i p^i / \Gamma^i)$, where p^{α} is the partial pressure of component α . Note that the formulation is sufficiently general to allow real gas EOS systems described by pressure given as a function of density and internal energy.

The solution procedure for the system is a higher-order Godunov method following that of [10] and detailed in [11]. This base integrator for the multifluid system is embedded within an implementation of adaptive mesh refinement (AMR). This methodology is based on the original work of [12] and later [13]. This is a means for managing a refined grid heirarchy composed of logically rectangular grid patches that allows for efficient increases in resolution by focusing the computational accuracy where errors are deemed large, or in regions where high resolution is required.

3. Problem Setup

The initial conditions for this problem are taken from [7]. The shock wave Mach number is 1.094 in air. In [14], a Raliegh scattering technique is used to measure the actual molar concentrations of Helium. It shows the initial profile is Gaussian in shape with a core concentration of approximately 80% Helium. The initial half radius is set as, $r_{half} = 0.2$, where it is defined as the distance where the Helium mole fraction is half the center value. The density ratio between air and Helium is $\rho_{Helium}/\rho_{air} = .138$.

The computational domain is 9cm wide by 27cm long with a base grid of 36 cells wide by 108 cells long. There are two levels of refinement with the first level is a factor of 4 more refined than the base grid and the second level is a factor of 8 more refined than the first level of refinement. This gives an effective grid resolution of $\Delta x_{\rm fine} = \Delta y_{\rm fine} = 0.0078125 {\rm cm}$. The shock is initialized in air a distance of $3r_{\rm half}$ from the jet. All boundaries are inflow/outflow type.

4. Results and Discussion

4.1. Flow Visualization

Figure 1 and 2 shows the evolution of the mole fraction of Helium and the development of the vorticity field, respectively. Figure 1(a) shows the initial conditions and the initial mixed layer. At 88μ sec, figure 1(b), there is a flattening of the jet and a generally weak deformation of its original circular structure. From here out to about 400μ sec, spanned in time by figures 1(c) through (f), there is stong deformation of the jet as it is inverted and transformed into the vortex pair. After 400μ sec, from figure 1(g) on, the flow evolves as a well defined vortex pair. Comparing these results with those given in [7], the agreement is exceptional.

Based on these observations and for ease and clarity of later exposition, we define three phases of development. The first is the weak deformation phase (WD) that occurs over short times out to about $90\mu\text{sec}$. The second is the strong deformation phase (SD) that occurs after WD out to about $400\mu\text{sec}$. The last is the relaxation/reorganization (RR) phase that occurs after SD.

In the vorticity time sequence, figure 2(a) shows the shock, still essentially planar, leaving positive (in the upper half plane) vorticity behind it. By 66μ sec, figure 2(b), the shock is well out of the jet core and two counter-sign circular vortex blobs are seen located in the jet core. During the WD phase, the vortex blobs remain essentially undistorted. This is due to the fact that vorticity induced velocity effects are very weak during the WD phase. Also, the blobs are deposited in the jet core over a timescale that scales with the shock passage time ($\sim 10\mu$ sec). During the SD phase however, there are large induced velocity effects that distort and stretch the vortex blobs into something like continuous finite thickness sheets. The underlying vortical field swirls these sheets so that they are reorganized, at later times, into a vortex pair. Note that during the SD phase, there is strong counter-sign vorticity appearing also oriented in sheets, ahead of the distorted blob. Although not shown here, this counter-sign production acts to modulate the overall circulation as it relaxes to a near constant value.

The development of this flow is seen to be determined by the interaction and evolution of the vortex blobs. Once they are present in the flow, they at first interact weakly (during WD), then interact strongly (during SD) as they change from blobs, to more point-like vortices (during RR). During the course of the transformation in the vorticity field, the Helium jet undergoes a similar transformation as shown above. From a modeling point of view, this flow could be accurately depicted by the development of a vortex blob pair in an essentially circular helium jet (slightly flattened on the upstream side in

reality). The subsequent development could be predicted by their interaction and reorganization/relaxation to point vortices.

References

- [1] R.D. Richtmyer 1960 Taylor instability in shock acceleration of compressible fluids. *Commun. Pure Appl. Maths* 23, 297-319.
- [2] E.E. Meshkov 1969 Instability of the interface of two gases accelerated by a shock wave. *Izv. Akad. Nauk. SSSR Mekh. Ahidk. Gaza* 4, 1151-157 [Russian: *Izv. Acad. Sci. USSR Fluid Dyn.* 4, 101-104.
- [3] Zaitsev, S.G., Lazareva, E.V., Chernukha, V.V., & Belyaev, V.M. 1985 Experimental investigation of the hydrodynamic instability of the interface between media of different density in an acceleration field. Translated from Teplofzika Vysokikh Temperatur 23, 535-541.
- [4] Brouillette, M. & Sturtevant, B. 1988 Shock induced Rayleigh-Taylor instability at a continuous interface. Abstract submitted to the *Intl Workshop on the Physics of Compressible Turbulent Mixing*, 24-27 October 1988, Princeton, New Jersey.
- [5] Haas, J.-F. & Sturtevant, B. 1987 Interaction of weak shock waves with cylindrical and spherical inhomogeneities. J. Fluid Mech., 181, 42-76.
- [6] Jacobs, J.W. 1992 Shock-induced mixing of a light-gas cylinder. J. Fluid Mech. 234, 629-649.
- [7] Jacobs, J.W. 1993 The dynamics of shock accelerated light and heavy gas cylinders. *Phys. Fluids* A 5, 2239-2247.
- [8] Colella, P., Ferguson, R.E., Glaz, H.M. 1995 Multifluid algorithms for eulerian finite difference methods. Manuscript.
- [9] Puckett, E.G. & Saltman, J.S. 1992 A 3-d adaptive mesh refinement algorithm for multimaterial gas dynamics. *Physica* D **60**, 84-104.
- [10] Colella, P. 1984 A direct eulerian muscl scheme for gas dynamics. SIAM J. Sci. Stat. Comput. 6, 104-117.
- [11] Greenough, J.A., Bell, J.B., Colella, P. 1995 An adaptive multifluid interfacecapturing method for compressible flows in complex geometry. AIAA Paper 95-1718.
- [12] Berger, M.J. & Oliger, J. 1984 Adaptive mesh refinement for hyperbolic partial differential equations. J. Comput. Phys. 53, 484-512.
- [13] Berger, M.J. & Colella P. 1989 Local adaptive mesh refinement for shock hydrodynamics. J. Comput. Physis., 82, 64-84.
- [14] Budzinski, J.M. 1992 Planar rayleigh scattering measurements of shock enhanced mixing. Doctoral dissertation. California Institute of Technology.

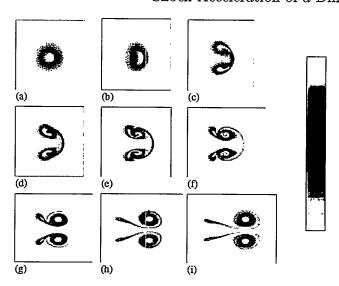


Figure 1. A sequence of the computed Helium mole fraction evolution. Note that the shock arrival time at the jet is approximately 40μ sec and all times reported are relative to the start of the calculation. (a) The inital jet, (b) 88μ sec, (c) 202μ sec, (d) 259μ sec, (e) 318μ sec, (f) 436μ sec, (g) 561μ sec, (h) 943μ sec, (i) 1182μ sec. The shock moves from left to right. The color bar to the right varies highest to lowest values from top to bottom.

