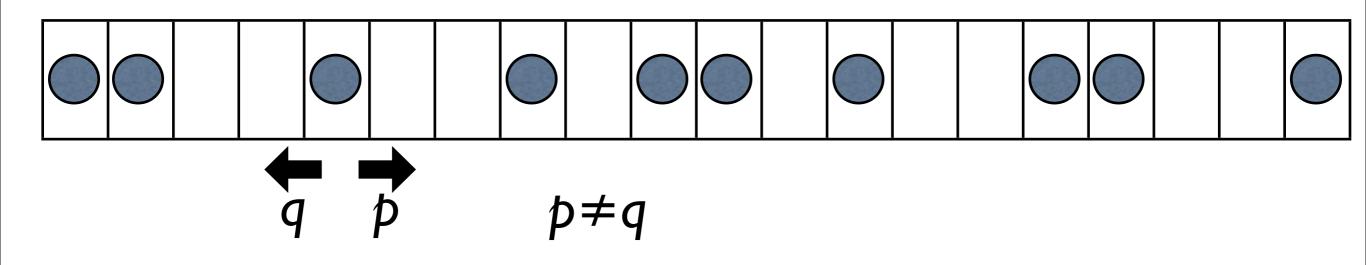
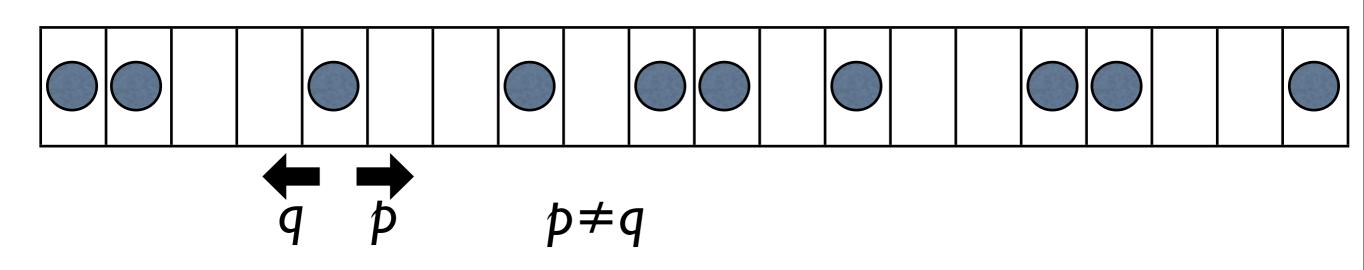
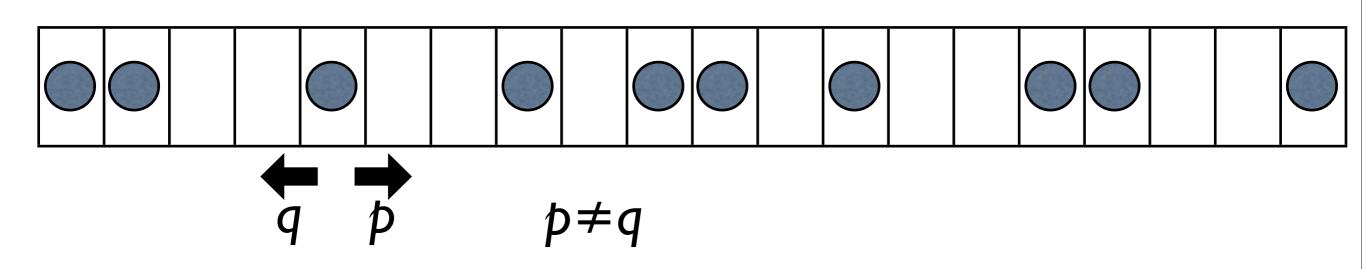
Asymmetric Simple Exclusion Process: Integrable Structure & Limit Theorems Northeast Probability Seminar November 2009

Craig Tracy & Harold Widom

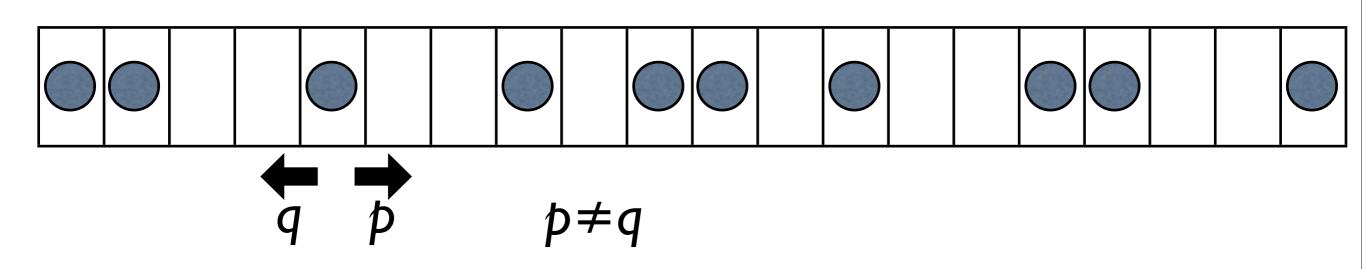




Each particle has an alarm clock - exponential distribution with parameter one



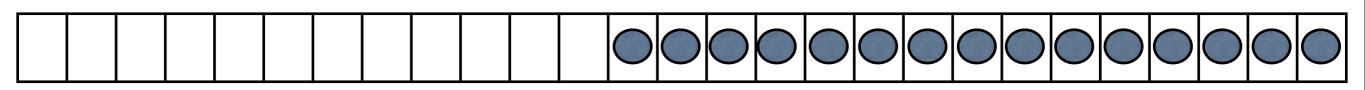
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- Each particle has an alarm clock -exponential distribution with parameter one
- When alarm rings particle jumps to right with probability p and to the left with probability q
- Jumps are suppressed if neighbor is occupied

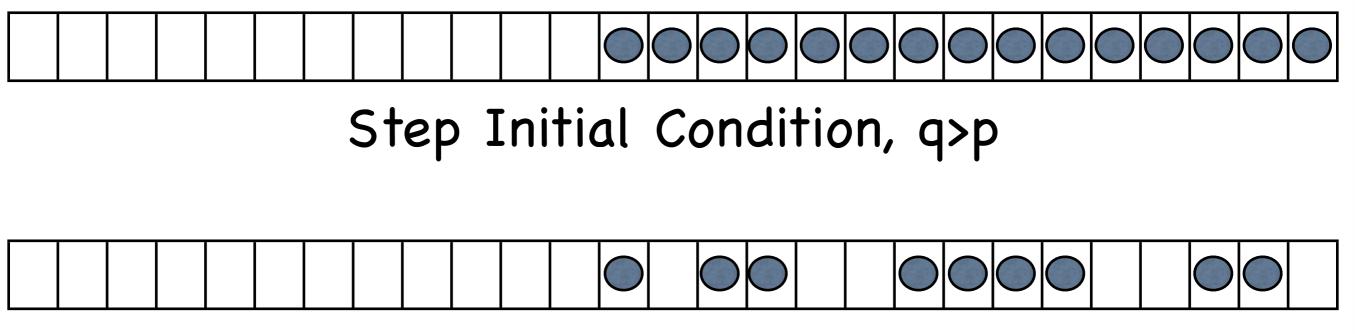
Initial Conditions



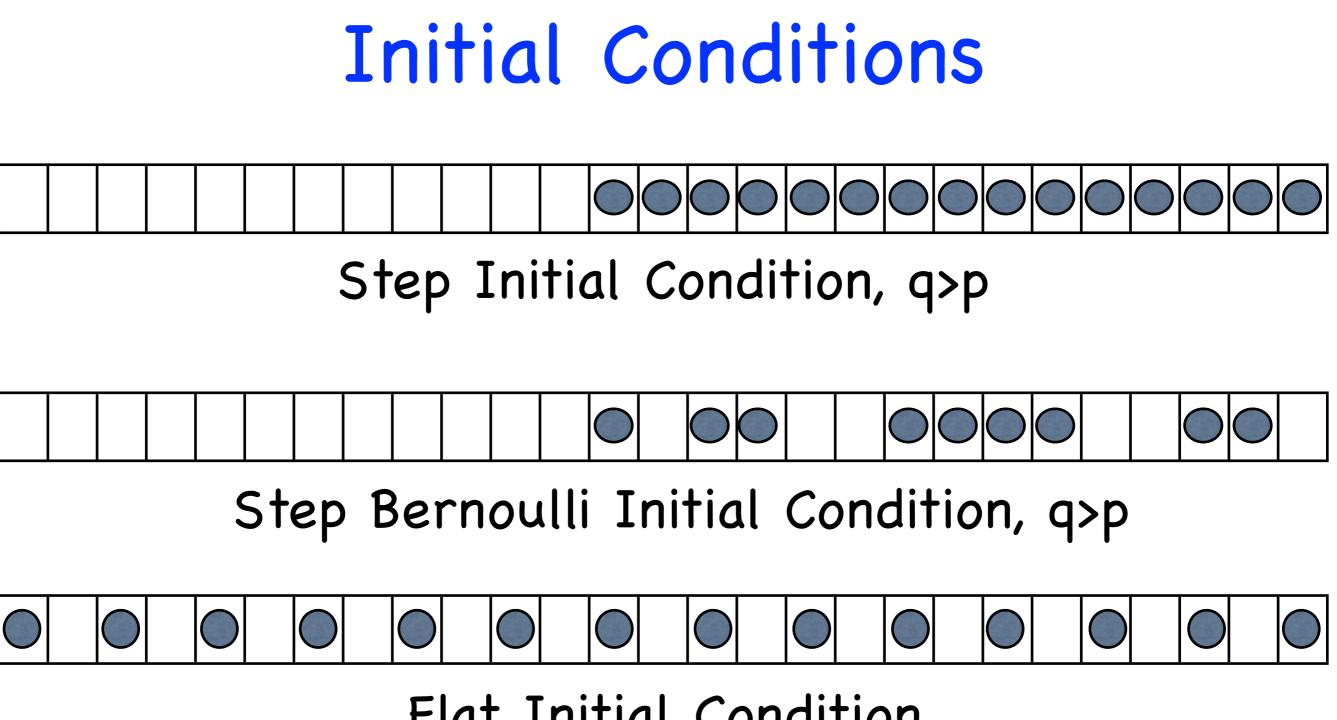


Step Initial Condition, q>p

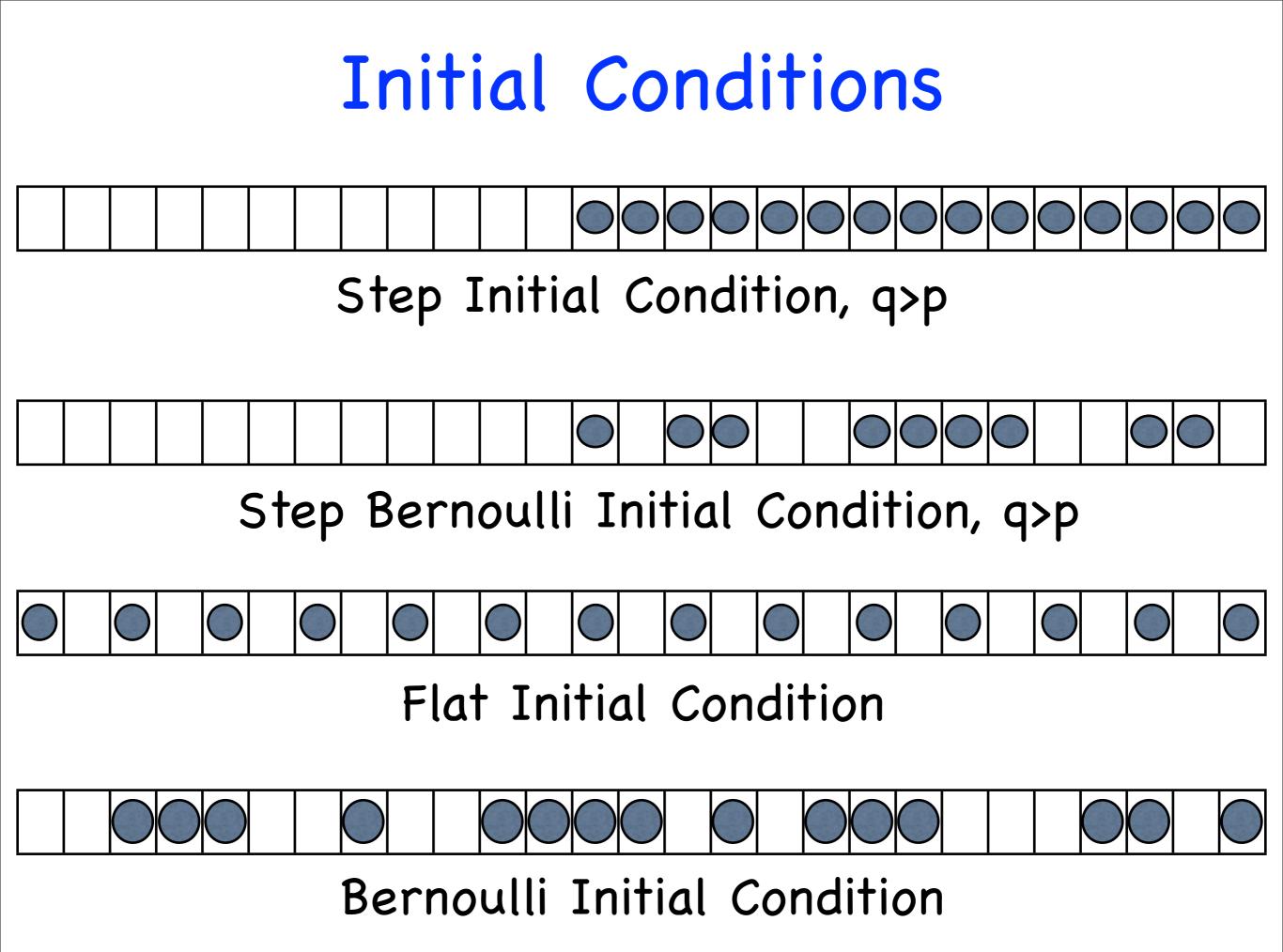




Step Bernoulli Initial Condition, q>p



Flat Initial Condition



Current Fluctuations

- Let I(x,t)= # of particles less than or equal to x at time t. (Definition OK for step type initial conditions.)
- I(x,t) is a random variable: What can be said about its long-time behavior? Can we find a central limit type theorem for I(x,t)?
- First some results for the special case of TASEP where particles are allowed to jump in only one direction

T(Totally)ASEP

- Determinantal Process: Correlations are expressed as determinants of kernel functions K(x,y).
- As a result many techniques from random matrix theory can be applied to TASEP.
- This determinantal structure goes back at least to G. Schütz (1997).
- First limit theorems: K. Johansson (2000), "Shape Fluctuations and Random Matrices" where TASEP was a limiting case of a certain Corner Growth Model.

Step initial condition xm(t)=position of mth particle from left at time t, xm(0)=m

Step initial condition x_m(t)=position of mth particle from left at time t, x_m(0)=m

Event: {I(x,t)=m}={x_m(t)≤x, x_{m+1}(t)>x}

Step initial condition x_m(t)=position of mth particle from left at time t, x_m(0)=m

Event: {I(x,t)=m}={x_m(t)≤x, x_{m+1}(t)>x}

From this and exclusion property:

 $Prob(I(x,t) \le m) = 1 - Prob(x_{m+1}(t) \le x)$

TASEP: Step Initial Condition

Theorem (K. Johansson):

$$\lim_{t \to \infty} \mathbb{P}\left(\frac{I([-vt], t) - a_1 t}{a_2 t^{1/3}} \le s\right) = 1 - F_2(-s)$$

where

$$a_1 = \frac{1}{4}(1-v)^2, \ a_2 = 2^{-4/3}(1-v^2)^{2/3}, \ 0 \le v < 1$$

and F₂ is a distribution function first arising in random matrix theory. It can be expressed in terms of a Painlevé II function or as a Fredholm determinant of the "Airy kernel".

Sunday, November 22, 2009

- Fluctuations are of order $t^{1/3}$ not $t^{1/2}$.
- This ¹/₃ exponent is related to the physicists'
 KPZ Universality (Kardar-Parisi-Zhang)
- The coefficient a_1 goes back to H. Rost (1981).
- Proof relies heavily on the RSK algorithm and the closely associated determinantal structure: "Discrete random martrix ensembles".

TASEP: Step Bernoulli, density p<1

Conjectured by M. Prähofer & H. Spohn (2002) & proved by G. Ben Arous & I. Corwin (2009):

 $\lim_{t \to \infty} \mathbb{P}\left(\frac{I([vt], t) - a_1 t}{a_2 t^{1/3}} \le s\right) = \begin{cases} 1 - F_2(-s), & -1 < v < 2\rho - 1, \\ 1 - F_1(-s)^2, & v = 2\rho - 1. \end{cases}$

where coefficients a_i as before $(v \rightarrow -v) \& F_1$ is another RMT distribution expressible in terms of the same Painlevé II function.

For v>2p-1 fluctuations are $t^{1/2}$ and Gaussian.

The transition: $F_2 \longrightarrow F_1^2 \longrightarrow G$ has appeared in other determinantal processes.

The Big Question

Do these results extend to ASEP

and more generally to a wider class of 1D exclusion processes ?

- With regards to the t^{1/3}: M. Balázs & T. Seppäläinen have proved this for Bernoulli (stationary) initial conditions for nearest-neighbor ASEP. Quastel & Valkó extend the result to finite-range asymmetric exclusion. Methods use coupling (BS) and comparison estimates (QV). Methods so far do not extend to limit theorems.
- H. Widom & CT show for (nearest neighbor) ASEP the two previous limit theorems remain exactly the same when time t is replaced by t/(q-p) thus proving KPZ Universality. Our methods start with ideas coming from Bethe Ansatz.

Main Steps in Proof of Limit Theorems for ASEP

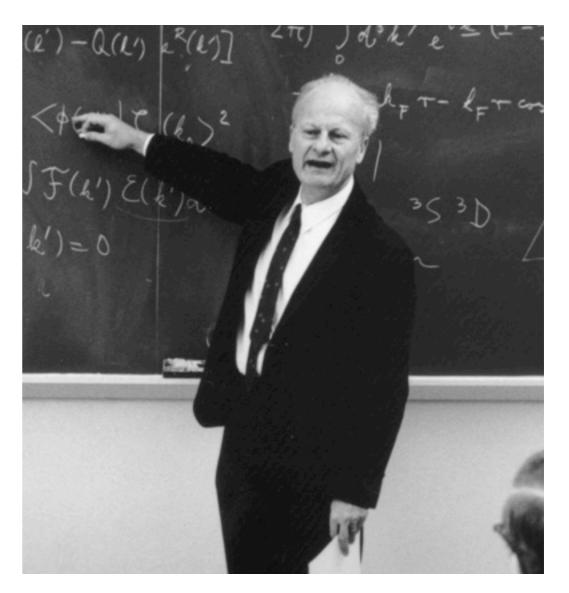
- Use ideas from Bethe Ansatz to find exact expression for the transition probability Y→X for finite N ASEP.
- Use some amazing combinatorial identities to compute marginal distributions P(x_m(t) ≤ x). It is then possible to take N → ∞. Marginal expressed as an infinite sum where kth term is a k-dimensional integral.
- Again certain identities permit series to be summed to one contour integral involving a Fredholm determinant.
- Asymptotic analysis of this Fredholm determinant.
 Kernel initially not of familiar RMT structure.

Integrable Structure of ASEP

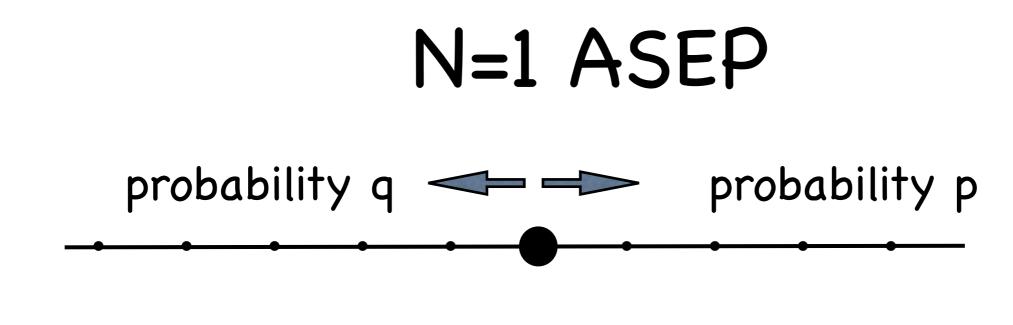
We solve the Kolmogorov forward equation ("master equation") for the transition probability $Y \rightarrow X$ for finite N ASEP:

 $P_Y(X;t)$

Main idea comes from the Bethe Ansatz (1931)



Hans Bethe in 1967



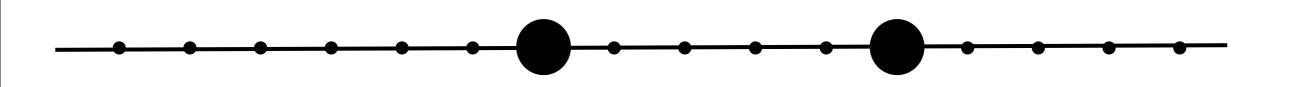
Let $P_Y(X;t)$ =probability $Y \rightarrow X$ at time t. Master equation:

$$\frac{dP}{dt} = p P(x - 1; t) + q P(x + 1; t) - P(x; t)$$

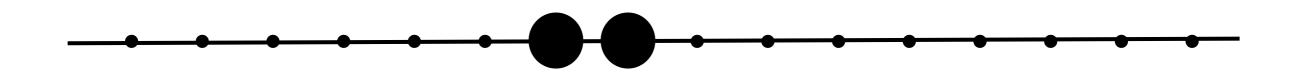
$$P_y(x;t) = \int_{\mathcal{C}} \xi^{x-y-1} e^{t\varepsilon(\xi)} d\xi$$

$$\varepsilon(\xi) = \frac{p}{\xi} + q\,\xi - 1$$

N=2 ASEP

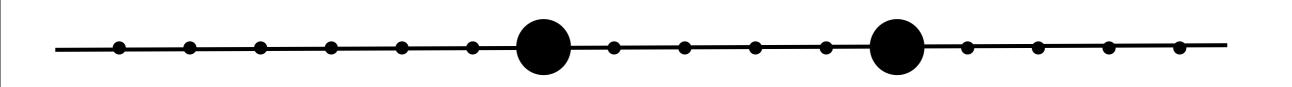


Master equation takes simple form for this configuration

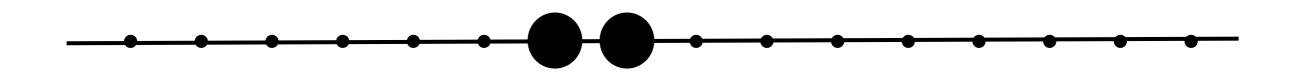


Master equation reflects exclusion for this configuration

N=2 ASEP



Master equation takes simple form for this configuration



Master equation reflects exclusion for this configuration

Impose boundary conditions for first equation so that if satisfied the second equation is automatically satisfied --- Bethe's Idea

N=2 Equations for
$$P(x_1, x_2; t)$$

 $x_2 > x_1 + 1$: Not neighbors

$$\frac{dP(x_1, x_2; t)}{dt} = pP(x_1 - 1, x_2; t) + qP(x_1 + 1, x_2; t) + pP(x_1, x_2 - 1; t) + qP(x_1, x_2 + 1; t) - 2P(x_1, x_2; t)$$

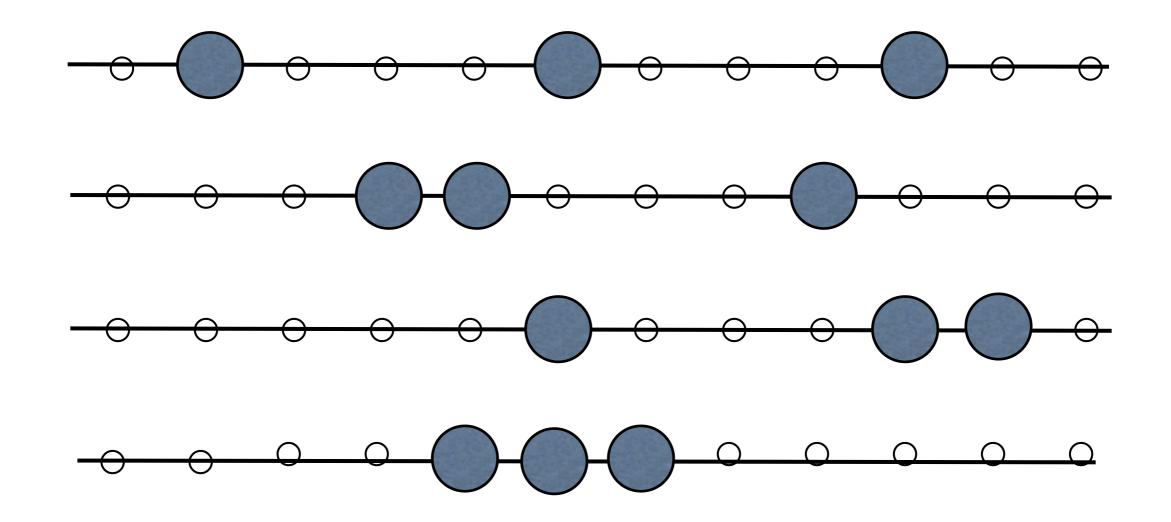
X₂=X₁+1: Neighbors $\frac{dP(x_1, x_2; t)}{dt} = pP(x_1 - 1, x_2; t) + qP(x_1, x_2 + 1; t) - P(x_1, x_2; t)$

Subtract:

 $qP(x_1 + 1, x_1 + 1; t) + pP(x_1, x_1; t) - P(x_1, x_1 + 1; t) = 0$

Important Point

New boundary conditions arise for N=3, 4,...



Last configuration requires new BC -automatically satisfied by 2-particle BC

Bethe Ansatz Solution of Master Equation

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For any $\xi_1, \dots, \xi_N \in \mathbb{C} \setminus \{0\}$ and any permutation σ a solution is $\prod \xi_{\sigma(j)}^{x_j} e^{t\varepsilon(\xi_j)}$

1

Bethe Ansatz Solution of Master Equation

For any $\xi_1, \dots, \xi_N \in \mathbb{C} \setminus \{0\}$ and any permutation σ a solution is $\prod_j \xi_{\sigma(j)}^{x_j} e^{t\varepsilon(\xi_j)}$

Can take linear combination or integral of a linear combination & have a solution:

$$\int \sum_{\sigma \in \mathcal{S}_N} F_{\sigma}(\xi) \prod_j \xi_{\sigma(j)}^{x_j} \prod_j e^{t\varepsilon(\xi_j)} d^N \xi$$

• Want BC to be satisfied. Bethe's idea: This can be applied pointwise to the integrand

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- Gives condition on coefficients with result (This part same as the C.N. Yang & C.P. Yang analysis of XXZ spin Hamiltonian.)

If
$$A_{\sigma} = \operatorname{sgn}(\sigma) \frac{\prod_{i < j} (p + q\xi_{\sigma(i)}\xi_{\sigma(j)} - \xi_{\sigma(i)})}{\prod_{i < j} (p + q\xi_i\xi_j - xi_i)}$$

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then solution to DE that satisfies BC is

$$\sum_{\sigma} \int A_{\sigma}(\xi) \prod_{i} \xi_{\sigma(i)}^{x_i - y_{\sigma(i)} - 1} e^{t\varepsilon(\xi_i)} d^N \xi$$

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- Must show the solution at X={x₁,...x_N} & t=0 reduces to $\delta_{X,Y}$ where Y={y₁,...,y_N} is the initial configuration of N particles.
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- Term for σ=id satisfies initial condition. Must show remaining N!-1 terms give zero at t=0.
 This is the new part of the problem.
- Have not yet specified the contours.

Theorem (TW): If $p \neq 0$ and r is small enough then

$$\mathbb{P}_Y(X;t) = \sum_{\sigma \in \mathcal{S}_N} \int_{\mathcal{C}_r^N} A_\sigma(\xi) \prod_i \xi_{\sigma(i)}^{x_i} \prod_i \left(\xi_i^{-y_i - 1} e^{\varepsilon(\xi_i) t} \right) d^N \xi.$$

where

$$A_{\sigma} = \operatorname{sgn} \sigma \frac{\prod_{i < j} (p + q\xi_{\sigma(i)}\xi_{\sigma(j)} - \xi_{\sigma(i)})}{\prod_{i < j} (p + q\xi_i\xi_j - \xi_i)}$$

and satisfies

$$\mathbb{P}_Y(X;0) = \delta_Y(X).$$

- There is no Ansatz in our work!
- Usual Bethe Ansatz calculates the spectrum of the operator. This leads to transcendental equations for the eigenvalues and issues of completeness of the eigenfunctions.
- We compute the semigroup directly. No spectral theory.

Marginal Distributions $P(x_m(t) \le x)$

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Case m=1:

Fix $x_1=x$, sum $P_Y(X;t)$ over allowed x_2 , x_3 , x_4 ,...

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Marginal Distributions $P(x_m(t) \le x)$

Case m=1:

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Can do this since contours are small: $|\xi_i| < 1$

Result is an expression involving N! terms. Use first miraculous identity to reduce sum to one term!

Here's the identity:

First Identity

$$\sum_{\sigma \in \mathcal{S}_N} \operatorname{sgn} \sigma \left(\prod_{i < j} (p + q\xi_{\sigma(i)}\xi_{\sigma(j)} - \xi_{\sigma(i)}) \right) \\ \times \frac{\xi_{\sigma(2)}\xi_{\sigma(3)}^2 \cdots \xi_{\sigma(N)}^{N-1}}{(1 - \xi_{\sigma(2)}\xi_{\sigma(3)} \cdots \xi_{\sigma(N)})(1 - \xi_{\sigma(3)} \cdots \xi_{\sigma(N)}) \cdots (1 - \xi_{\sigma(N)})} \right) \\ = p^{N(N-1)/2} \frac{(1 - \xi_1 \cdots \xi_N) \prod_{i < j} (\xi_j - \xi_i)}{\prod_i (1 - \xi_i)}$$

● Using this identity we get for m=1 an expression for P(x1(t)≤x) as a single Ndimensional integral with a product integrand. This expression is for finite-N ASEP

$$I(x,Y,\xi) = \prod_{i < j} \frac{\xi_j - \xi_i}{p + q\xi_i\xi_j - \xi_i} \frac{1 - \xi_1 \cdots \xi_N}{(1 - \xi_1) \cdots (1 - \xi_N)}$$
$$\prod_i \left(\xi_i^{x - y_i - 1} e^{\varepsilon(\xi_i)t} \right)$$

 $Prob(x_1(t)=x) =$ $p^{N(N-1)/2} \int_{\mathcal{C}_r} \cdots \int_{\mathcal{C}_r} I(x, Y, \xi) d\xi_1 \cdots d\xi_N$ $(p \neq 0)$

$$\begin{aligned} \mathsf{Prob}(\mathsf{x}_1(\mathsf{t})=\mathsf{x}) &= \\ p^{N(N-1)/2} \int_{\mathcal{C}_r} \cdots \int_{\mathcal{C}_r} I(x, Y, \xi) \, d\xi_1 \cdots d\xi_N \\ (p \neq 0) \end{aligned}$$

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- Sum of N! integrals reduced to one integral
- Form is not so useful to take $N\!\rightarrow\!\infty$
- We now expand contour outwards -- only residues that contribute come from $\xi_i=1$.
- Can take $N \rightarrow \infty$ in resulting expression to obtain

$$\sigma(S) := \sum_{i \in S} i$$

$$\mathbb{P}(x_1(t) = x) = \sum_{S} \frac{p^{\sigma(S) - |S|}}{q^{\sigma(S) - |S|(|S| + 1)/2}} \times \int_{\mathcal{C}_R} \cdots \int_{\mathcal{C}_R} I(x, Y_S, \xi) d^{|S|} \xi$$

The sum is over all nonempty subsets of \mathbb{Z}^+ (finite)

When p=0 only one term is nonzero, S={1}.

To go beyond the left-most particle, m=1, there are new complications

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• This involves finding a new identity

Second Identity

S ranges over subsets of $\{1, 2, \ldots, N\}$

$$\sum_{|S|=m} \prod_{i \in S, j \in S^c} \frac{p + q\xi_i \xi_j - \xi_i}{\xi_j - \xi_i} \cdot (1 - \prod_{j \in S^c} \xi_j)$$

$$= q^m \begin{bmatrix} N \\ m \end{bmatrix} (1 - \prod_{j=1}^N \xi_j).$$

$$[N] = \frac{p^{N} - q^{N}}{p - q}, \qquad [N]! = [N] [N - 1] \cdots [1],$$

$$N] \qquad [N]! \qquad (n + 1) \cdots (n + 1) \cdots (n + 1)$$

$$\begin{bmatrix} n \end{bmatrix} = \frac{[n]!}{[m]! [N-m]!}, \qquad (q-\text{binomial coefficient}),$$

$$\mathbb{P}\left(x_{m}(t) \leq x\right) = (-1)^{m} \sum_{k \geq m} \frac{1}{k!} \begin{bmatrix} k-1\\ k-m \end{bmatrix}_{\tau} p^{(k-m)(k-m+1)/2} q^{km+(k-m)(k+m-1)/2} \\ \times \int_{\mathcal{C}_{R}} \cdots \int_{\mathcal{C}_{R}} \prod_{i \neq j} \frac{\xi_{j} - \xi_{i}}{p + q\xi_{i}\xi_{j} - \xi_{i}} \prod_{i} \frac{1}{(1-\xi_{i})(q\xi_{i}-p)} \\ \times \prod_{i} \left(\xi_{i}^{x} e^{\varepsilon(\xi_{i})t}\right) d\xi_{1} \cdots d\xi_{k}$$

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$$\times \prod_{i} \left(\xi_{i}^{x} e^{\varepsilon(\xi_{i})t}\right) d\xi_{1} \cdots d\xi_{k}$$

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- For p=0 only k=m term is nonzero
- Recognize double product as a determinant whose entries are a kernel, i.e. $K(\xi_i,\xi_j)$
- Result can then be expressed as a contour integral whose integrand is a Fredholm determinant.

Fredholm determinant

- Let K(x,y) be a kernel function
- Fredholm expansion of det(I- λ K):

$$\frac{(-1)^n}{n!} \int \cdots \int \det \left(K(\xi_i, \xi_j)_{1 \le i,j \le n} \ d\xi_1 \cdots d\xi_n = \int_{\mathcal{C}} \det \left(I - \lambda K \right) \ \frac{d\lambda}{\lambda^{n+1}}$$

•Can then do sum over k (q-Binomial theorem):

Final expression for mth particle distribution fn. Step initial condition

Set $\gamma = q - p > 0$, $\tau = p/q < 1$ and define an integral operator K on the circle C_R , $R \gg 1$

$$K(\xi,\xi') = q \, \frac{\xi^x e^{\varepsilon(\xi)t}}{p + q\xi\xi' - \xi}$$

then

$$\mathbb{P}\left(x_m(t/\gamma) \le x\right) = \int \frac{\det\left(I - \lambda K\right)}{\prod_{k=0}^{m-1} \left(1 - \lambda \tau^k\right)} \frac{d\lambda}{\lambda}$$

where the contour encloses all the singularities at $\lambda=0, \tau^{-k}$, k=0,...,m-1.

Final expression for mth particle distribution function Step Bernoulli initial condition

The expression for the marginal $P(x_m(t) \le x)$

is the same once we replace K by the new K:

$$K(\xi,\xi') = q \frac{\xi^{x} e^{\varepsilon(\xi)t}}{p + q\xi\xi' - \xi} \frac{\rho(\xi - \tau)}{\xi - 1 + \rho(1 - \tau)}$$

$$\uparrow$$
new term

Asymptotic analysis

We now transform the operator K so that we can perform a steepest descent analysis.

Recall that the generic behavior for the coalescence of two saddle points leads to the Airy function Ai(x)



George Airy

$$\begin{split} \xi &\longrightarrow \quad \frac{1-\tau\eta}{1-\eta}, \quad \tau = \frac{p}{q} < 1, \\ K(\xi,\xi') &\longrightarrow \quad K_2(\eta,\eta') = \frac{\varphi(\eta')}{\eta'-\tau\eta} \\ \varphi(\eta) &= \quad \left(\frac{1-\tau\eta}{1-\eta}\right)^x \ e^{\left[\frac{1}{1-\eta} - \frac{1}{1-\tau\eta}\right] \mathbf{t}} \end{split}$$

Introduce:
$$K_1(\eta, \eta') = \frac{\varphi(\tau \eta)}{\eta' - \tau \eta}$$

.

Two Preliminary Results

Two Preliminary Results Proposition:

Let Γ be any closed curve going around $\eta=1$ once counterclockwise with $\eta=1/\tau$ on the outside. Then the Fredholm determinant of K (ξ,ξ') acting on C_R has the same Fredholm determinant as $K_1(\eta,\eta')-K_2(\eta,\eta')$ acting on Γ .

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Proposition:

Suppose the contour Γ is star-shaped with respect to $\eta=0$. Then the Fredholm determinant of K₁ acting on Γ is equal to

$$\prod_{k=0}^{\infty} (1 - \lambda \tau^k)$$

Denote by R the resolvent kernel of K₁

Factor determinant: det(I-λ K)=det(I-λ K₁) det(I+K₂(I+R))

Set $\lambda = \tau^{-m} \mu$ so formula for distr. fn becomes

$$\int \prod_{k=0}^{\infty} (1 - \mu \tau^k) \det \left(I + \tau^{-m} \mu K_2 (I + R) \right) \frac{d\mu}{\mu}$$

μ runs over a circle of radius > τ

By a perturbative expansion of R, followed by a deformation of operators, we show

$$\det (I + \lambda K_2(I + R)) = \det (I + \mu J)$$

$$J(\eta, \eta') = \int \frac{\varphi_{\infty}(\zeta)}{\varphi_{\infty}(\eta')} \frac{\zeta^m}{(\eta')^{m+1}} \frac{f(\mu, \zeta/\eta')}{\zeta - \eta} d\zeta$$

$$\varphi_{\infty}(\eta) = (1 - \eta)^{-x} e^{\frac{\eta t}{1 - \eta}}$$

$$f(\mu, z) = \sum_{k=-\infty}^{\infty} \frac{\tau^k}{1 - \tau^k \mu} z^k$$

The kernel $J(\eta, \eta')$, which acts on a circle centered at 0 with radius less than τ , is analyzed by the steepest descent method. Note: m now appears inside the kernel!

Main Result: Step I.C.

We set

$$\sigma = \frac{m}{t}, c_1 = -1 + 2\sqrt{\sigma}, c_2 = \sigma^{-1/6} (1 - \sqrt{\sigma})^{2/3}, \gamma = q - p$$

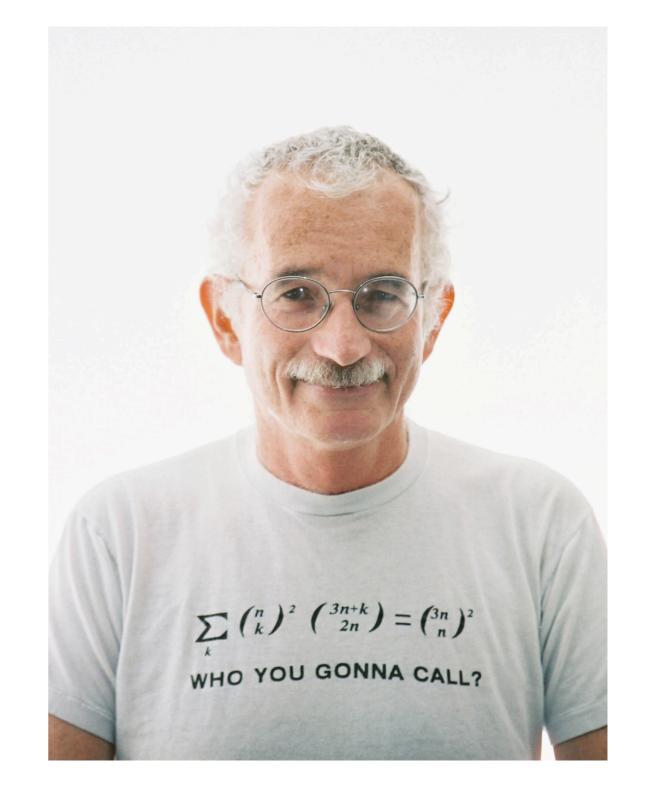
Theorem (TW). When $0 \le p < q$ we have

$$\lim_{t \to \infty} \mathbb{P}\left(\frac{x_m(t/\gamma) - c_1 t}{c_2 t^{1/3}} \le s\right) = F_2(s)$$

Theorem also has a current fluctuation formulation

Thanks to A. Schilling & D. Zeilberger for advice with the combinatorial identities





Thank you for your attention