

Lecture #20

The Riemann Tensor

We defined the Riemann tensor in terms of the operator, $u, v \in \mathcal{X}M$

$$R(u, v) : \mathcal{X}M \xrightarrow{\quad} \mathcal{X}M$$
$$w \mapsto \left(\nabla_u \nabla_v - \nabla_v \nabla_u - \nabla_{[u, v]} \right) w$$

and defined a $(1, 3)$ -tensor R by

$$R(w, \alpha, u, v) = \langle \alpha, R(u, v) w \rangle.$$

Components of the Riemann tensor

$$R(\partial_i, dx^k, \partial_m, \partial_n) = R_i{}^k{}_{mn}, \quad \nabla_m \equiv \nabla_{\partial_m}$$

$$= \langle dx^k, \nabla_m (\nabla_n \partial_i) - (m \leftrightarrow n) \rangle = \langle dx^k, \nabla_m (\underbrace{\Gamma_{ni}^j}_{\text{functions}} \partial_j) - m \leftrightarrow n \rangle$$

$$= \langle dx^k, \partial_m \Gamma_{ni}^j \partial_j + \Gamma_{ni}^j \Gamma_{mj}^l \partial_l - m \leftrightarrow n \rangle$$

$$= \boxed{\partial_m \Gamma_{ni}^k - \partial_n \Gamma_{mi}^k + \Gamma_{ni}^j \Gamma_{mj}^k - \Gamma_{mi}^j \Gamma_{nj}^k = R_i{}^k{}_{mn}}$$

Symmetries of the Riemann Tensor

From the coordinate expression, assuming vanishing torsion, it is not difficult to verify $R_{ijmn} = g_{jl} R_i{}^l{}_{mn}$ obeys

$$(i) \quad R_{ijmn} = R_{jimn}$$

$$(ii) \quad R_{ijmn} = R_{mnij}$$

$$(iii) \quad R_{ijmn} + R_{mjn i} + R_{jmin} = 0$$

(Bianchi identity of the 1st kind).

If $\dim M = n$ an antisymmetric matrix has $\frac{n(n-1)}{2}$ components so (i) & (ii) imply R_{ijmn} has

$$\frac{1}{2} \left(\frac{n(n-1)}{2} \right) \cdot \left(\frac{n(n-1)}{2} + 1 \right) = \frac{1}{2} n(n-1)(n^2 - n + 2)$$

components. Using (i) & (ii)

$R_{ijmn} + R_{mjn i} + R_{jmin}$ is totally antisymmetric

This gives $\binom{n}{4}$ additional constraints and leaves

$$\frac{1}{12} n^2 (n^2 - 1) \quad \text{components}$$

0	in 1 dim
1	in 2 dim
6	in 3 dim
20	in 4 dim.

2nd Bianchi Identity

Notice that in a coordinate basis $\{u, v, w, \dots\}$

$$R(u, v) = \nabla_u \nabla_v - \nabla_v \nabla_u$$

$$\Rightarrow \nabla_w R(u, v) + \text{cyclic}$$

$$= \nabla_w \nabla_u \nabla_v - \nabla_w \nabla_v \nabla_u$$

$$+ \nabla_u \nabla_v \nabla_w - \nabla_v \nabla_u \nabla_w$$

$$+ \nabla_v \nabla_w \nabla_u - \nabla_u \nabla_w \nabla_v$$

$$= 0$$

[In fact true in any basis]
but there are better ways
to prove it...

i.e.:

$$\nabla_k R_{ijmn} + \nabla_i R_{jkmn} + \nabla_j R_{kimn} = 0$$

Constant curvature manifolds

$$R(g^{-1}(\beta), \alpha, u, v) \propto \alpha(u)\beta(v) - \alpha(v)\beta(u)$$

$$\text{This says } R^{ij}_{mn} \propto \delta^i_m \delta^j_n - \delta^i_n \delta^j_m$$

Example The sphere with its canonical metric.

Irreducible decomposition of the Riemann Tensor

Let h, k be $\binom{0}{2}$ -tensors

$$\begin{aligned}(h \tilde{\otimes} k)(v_1, v_2, v_3, v_4) &= h(v_1, v_3)k(v_2, v_4) \\ &\quad - (1 \leftrightarrow 2) - (3 \leftrightarrow 4) \\ &\quad + (1 \leftrightarrow 2, 3 \leftrightarrow 4)\end{aligned}$$

This is sometimes called the Kulkarni Nomizu product and produces a tensor with the symmetries of Riemann. Then the $\binom{0}{4}$ -Riemann-tensor R

$$R = W + \frac{1}{n-2} (\text{Ric} - \frac{1}{n} s g) \tilde{\otimes} g + \frac{1}{2n(n-1)} s g \otimes g$$

Where W and Ric are trace-free $\binom{0}{4}$ and $\binom{0}{2}$ tensors. s is a scalar

This formula is perhaps easier to use in components:

$$R_{ijkl} - W_{ijkl} = g_{ik} P_{jl} - g_{jk} P_{il} - g_{il} P_{jk} + g_{jl} P_{ik}$$

where

W_{ijkl} is the Weyl tensor.

which has the symmetries of Riemann but

$$W^i{}_{jil} = 0$$

$P_{jl} = P_{lj}$ is the Schouten tensor

Notice $R - W = g \otimes P$

It is built from the Ricci and scalar curvatures

$$R_{ij} = R^k{}_{ijk}, \quad s = R^i{}_i$$

via

$$P_{ij} = \frac{1}{n-2} \left(R_{ij} - \frac{1}{2} \frac{1}{n-1} g_{ij} s \right)$$

* The Weyl Tensor is the obstruction to conformal flatness when $n \geq 4$.

i.e. We call a manifold flat when the Riemann tensor vanishes.

If $ds^2 = \Omega^2 ds_0^2$ where ds_0^2 is a flat metric (the Riemann tensor for its Levi-Civita connexion vanishes) we say that (M, ds^2) is conformally flat.

* In 3-dimensions the Weyl tensor vanishes identically (and instead the "covariant curl" of the P-tensor - called the Cotton tensor - measures conformal flatness).

- * There always exists a conformally equivalent metric such that the scalar curvature s is constant (Yamabe conjecture - now proven).

For a sphere it measures the inverse radius².

- * The Ricci tensor was employed by Einstein to write equations for the space time metric. In the absence of matter, he postulated

$$\text{Ric} = 0$$

(Such manifolds are called Ricci flat and include the famous Schwarzschild solution).

Review Exercises

- ① Compute the difference between the Riemann tensor for a connexion with torsion and for the Levi-Civita connexion. Which symmetries of the Riemann tensor fail?
- ② Compute the Riemann curvature for the sphere S^2 with canonical metric
$$ds^2 = d\theta^2 + \sin^2\theta d\phi^2.$$