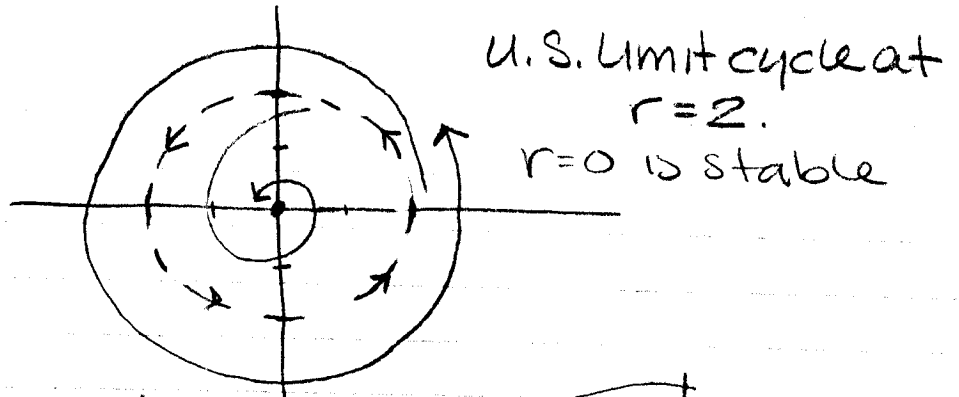


7.1.1 $\dot{r} = r^3 - 4r, \dot{\theta} = 1$

r -null: $\dot{r} = r(r^2 - 4) = 0 \Rightarrow r = 0, 2 \Rightarrow$ f.p. at $r = 0$

θ -null: DNE



7.1.2 $\dot{r} = r(1-r^2)(9-r^2), \dot{\theta} = 1$

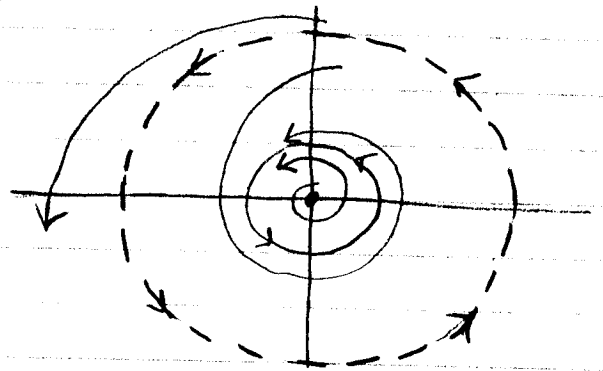
Limit cycles at $r = 1$ and 3

Fixed point at $r = 0$

$r = 1$ is a stable L.C.

$r = 3$ is an unstable L.C.

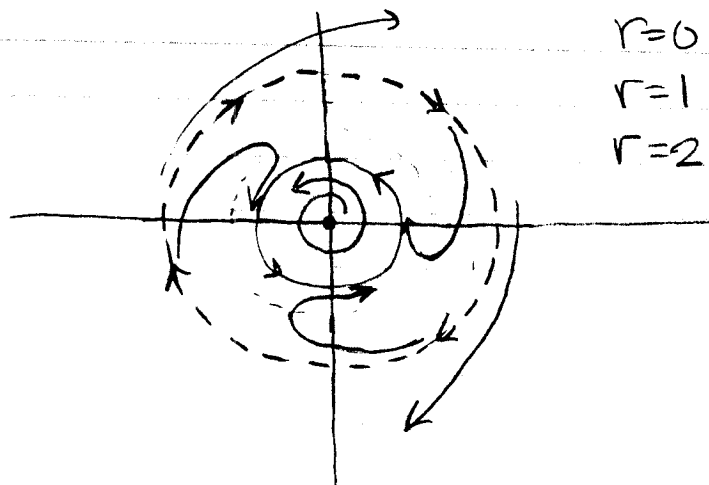
$r = 0$ is u.s.



7.1.3 $\dot{r} = r(1-r^2)(4-r^2), \dot{\theta} = 2-r^2$

$\dot{r} = 0$ at $r = 0, 1, 2$ } Fixed pt. at $r = 0$

$\dot{\theta} = 0$ at $r = \sqrt{2}$

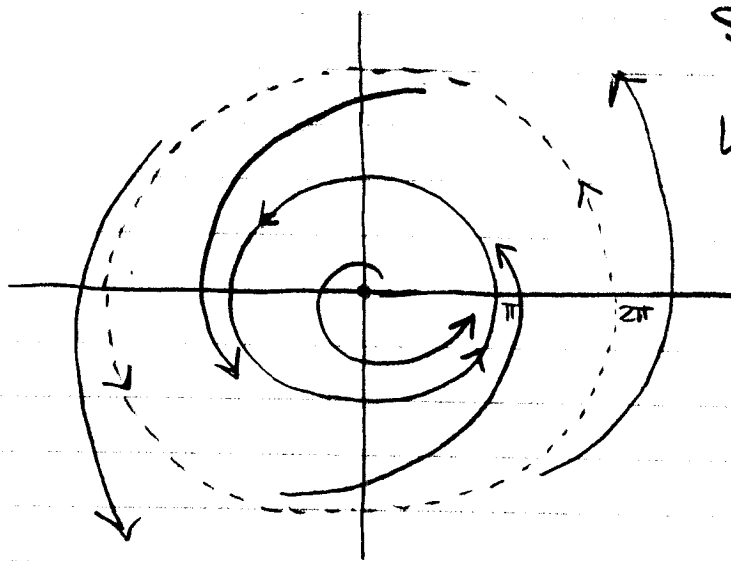


$r = 0$ is unstable
 $r = 1$ is a stable L.C.
 $r = 2$ is an unstable L.C.

7.1.4

$$\dot{r} = r \sin r, \dot{\theta} = 1$$

$\dot{r} = 0 \Rightarrow r = 0$ or $n\pi, n \in \mathbb{N}$ } $\Rightarrow r = 0$ is the only f.p.
 $\dot{\theta} \neq 0$



Stable limit cycles
when $r = (2n+1)\pi$;
U.S. limit cycles
when $r = 2n\pi$
($n \in \mathbb{N}$).

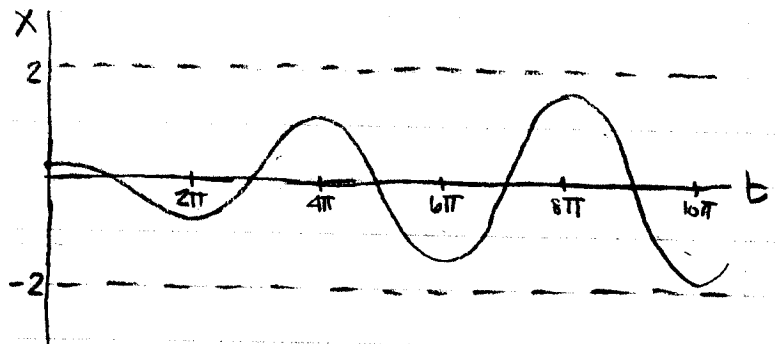
7.1.7 $\dot{r} = r(4-r^2), \dot{\theta} = 1$

$r=0$ is an unstable f.p. and $r=2$ is a stable L.c.
Period of cycles is 2π

\therefore If $x(0) = -1$ and $y(0) = 0$, then $x(t)$ will look

Something like

(the amplitude of oscillations will approach 2)



$$7.1.8 \quad \ddot{x} + a\dot{x}(x^2 + \dot{x}^2 - 1) + x = 0$$

$$a) \quad \begin{cases} \dot{x} = y \\ \dot{y} = -ay(x^2 + y^2 - 1) - x \end{cases}$$

fixed point. $(0, 0)$

Jacobian

$$A = \begin{bmatrix} 0 & 1 \\ -a2xy - 1 & -ax^2 - 3ay^2 + a \end{bmatrix}$$

$$A_{(0,0)} = \begin{bmatrix} 0 & 1 \\ -1 & a \end{bmatrix}$$

$$\Rightarrow \begin{aligned} -\lambda(a - \lambda) + 1 &= 0 \\ \lambda^2 - a\lambda + 1 &= 0 \end{aligned}$$

$$\Rightarrow \lambda = \frac{a \pm \sqrt{a^2 - 4}}{2}$$

\Rightarrow if $0 < a < 2 \Rightarrow (0, 0)$ is unstable spiral

if $a > 2 \Rightarrow (0, 0)$ unstable node

b) Convert to polar coordinates

$$\begin{aligned}r\dot{r} &= x\dot{x} + y\dot{y} \\ &= xy + y(-ay(x^2+y^2-1)-x) \\ &= \cancel{xy} - ay^2(x^2+y^2-1) - \cancel{xy} \\ &= -ay^2(x^2+y^2-1) \\ &= -a(r\sin\alpha)^2(r^2-1)\end{aligned}$$

$$\Rightarrow \dot{r} = -ar(r^2-1)\sin^2\alpha$$

$$\begin{aligned}r^2\dot{\alpha} &= \underline{x\dot{y} - y\dot{x}} \\ &= x(-ay^2(x^2+y^2-1)-x) - y(y) \\ &= -axy^2(x^2+y^2-1) - x^2 - y^2 \\ &= -axy(x^2+y^2-1) - (x^2+y^2) \\ &= -ar^2\cos\alpha\sin\alpha(r^2-1) - r^2\end{aligned}$$

$$\Rightarrow \dot{\alpha} = -a\cos\alpha\sin\alpha(r^2-1) - 1$$

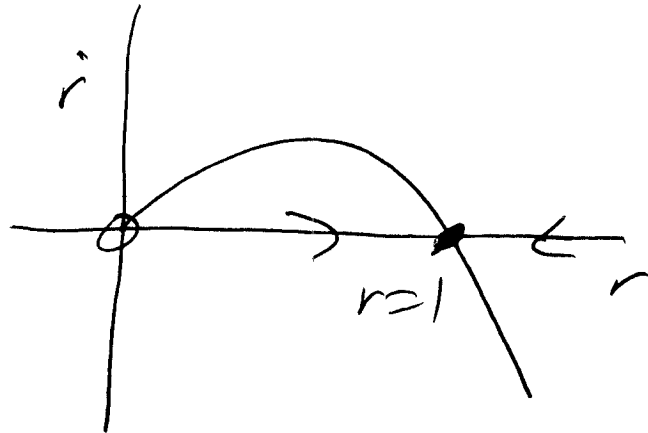
$$\Rightarrow \begin{cases} \dot{r} = -ar(r^2-1)\sin^2\alpha \\ \dot{\alpha} = -a\cos\alpha\sin\alpha(r^2-1) - 1 \end{cases}$$

Note that $\sin^2 \alpha > 0$

So, if $\sin^2 \alpha \neq 0$

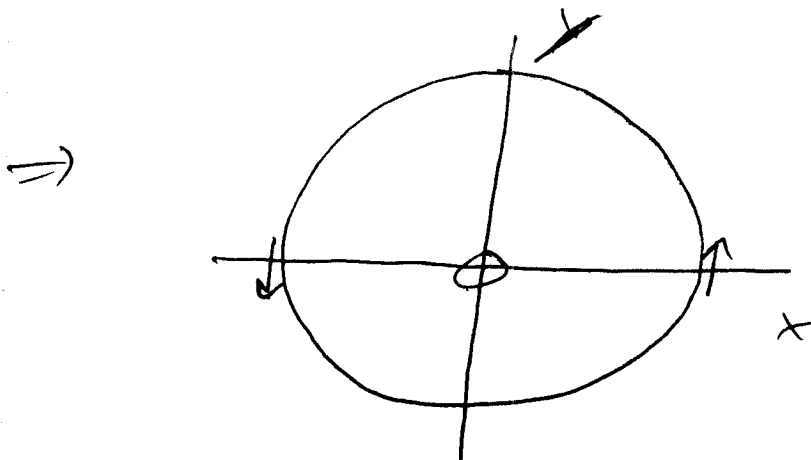
$\Rightarrow \dot{r} = 0$, so at 0 or $\pm\pi$, trajectories are strictly vertical

So, we can plot $-r(r^2-1)\tilde{\alpha}$ (absorb $\sin^2 \alpha$ into $\tilde{\alpha}$)



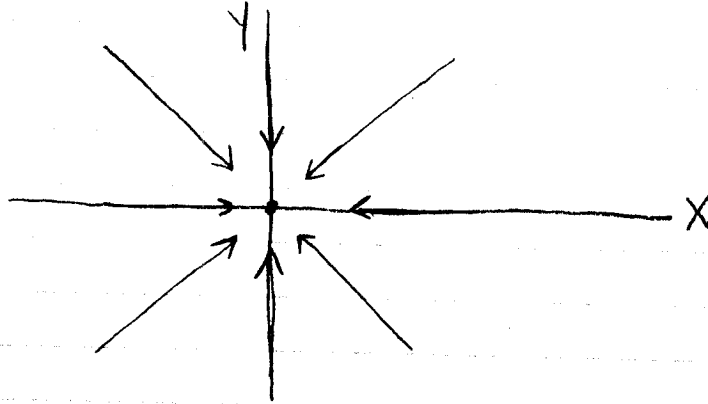
c) \Rightarrow existence of a stable limit cycle at $r=1$ and at $r=1$

$$\dot{\phi} = -1$$

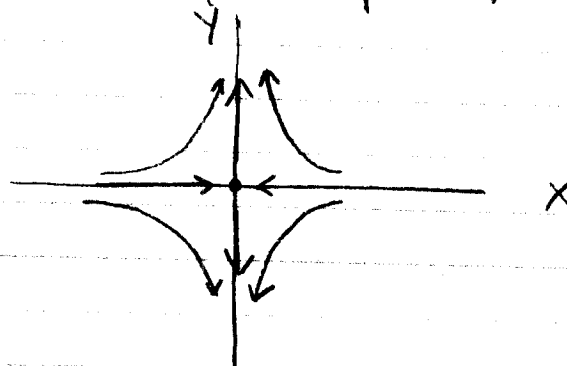


$$T = 2\pi$$
$$\text{amplitude} = 1$$

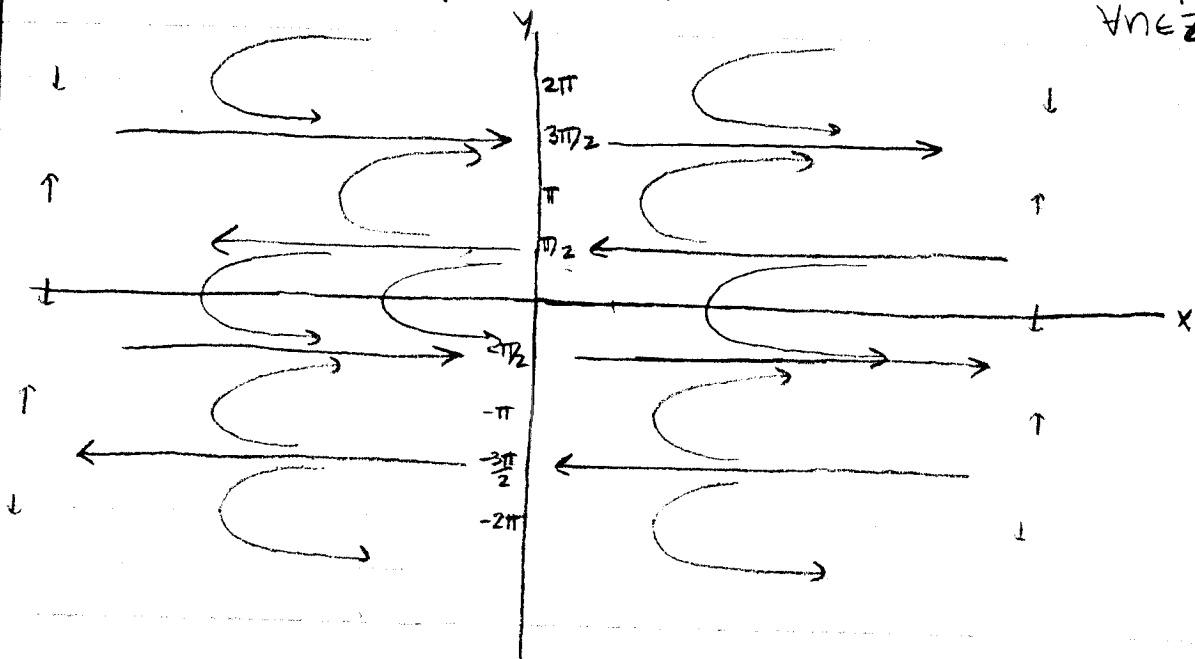
7.2.1 $V = x^2 + y^2 \Rightarrow \dot{x} = -2x, \dot{y} = -2y$
 $\Rightarrow (0,0)$ is the only fixed point, and it's stable



7.2.2 $V = x^2 - y^2 \Rightarrow \dot{x} = -2x, \dot{y} = 2y$
 $\Rightarrow (0,0)$ is the only fixed point, and it's a saddle.



7.2.3 $V = e^x \sin y \Rightarrow \dot{x} = -e^x \sin y, \dot{y} = -e^x \cos y$
 \Rightarrow no fixed points $y = 3\pi/2 + 2n\pi$ is a stable asymptote $\forall n \in \mathbb{Z}$



$$7.2.6 \ a) \ \begin{aligned} \dot{x} &= y^2 + y \cos x \\ \dot{y} &= 2xy + \sin x \end{aligned}$$

$$\begin{aligned} \dot{x} &= -dV/dx \\ \dot{y} &= -dV/dy \end{aligned}$$

$$\int -\dot{x} \, dx = -xy^2 - y \sin x$$

$$\int -\dot{y} \, dy = -xy^2 - y \sin x$$

Thus,

$$V(x, y) = -xy^2 - y \sin x$$

$$b) \ \begin{aligned} \dot{x} &= 3x^2 - 1 - e^{xy} \\ \dot{y} &= -2xe^{xy} \end{aligned}$$

$$\int -\dot{x} \, dx = -x^3 + x + xe^{xy}$$

$$\int -\dot{y} \, dy = xe^{xy}$$

Thus,

$$V(x, y) = -x^3 + x + xe^{xy}$$

7.2.10 $\dot{x} = y - x^3$, $\dot{y} = -x - y^3$, $V = ax^2 + by^2$

Show no closed orbits.

We need $V > 0$, $V(x^*, y^*) = 0$, $\dot{V} < 0$ along trajectories:

• $a, b > 0 \Rightarrow V > 0$

• $V(0, 0) = 0$

• $\dot{V} = 2ax\dot{x} + 2by\dot{y} = 2ax(y - x^3) + 2by(-x - y^3)$
 $= 2axy - 2ax^4 - 2bxy - 2by^4 < 0$ if $a = b$

• $\therefore V(x, y) = a(x^2 + y^2)$ is a Liapunov Fun $\forall a > 0$.
 \Rightarrow the system has no closed orbits.

? 2.11 $V = ax^2 + 2bxy + cy^2$ is positive definite if and only if $a > 0$ and $ac - b^2 > 0$.

Assume $a > 0$ & $ac - b^2 > 0$

$$V = ax^2 + 2bxy + cy^2$$

$$= a\left(x^2 + 2\frac{b}{a}xy\right) + cy^2$$

$$= a\left(x^2 + 2\frac{b}{a}xy + \left(\frac{b}{a}y\right)^2 - \left(\frac{b}{a}y\right)^2\right) + cy^2$$

$$= a\left(x + \frac{b}{a}y\right)^2 - a\left(\frac{b}{a}y\right)^2 + cy^2$$

$$= a\left(x + \frac{b}{a}y\right)^2 + y^2\left(c - \frac{b^2}{a}\right)$$

but $\left(x + \frac{b}{a}y\right)^2 > 0$
 $ac - b^2 > 0 \Rightarrow \left(c - \frac{b^2}{a}\right) > 0$
 $a > 0$
 $y^2 > 0$

Thus,

$$V > 0 \text{ . (positive definite)}$$

Next, suppose V is positive & finite
then

$$V = a\left(x + \frac{b}{a}y\right)^2 + y^2\left(ac - \frac{b^2}{a}\right) > 0$$

so, we need

$$a > 0$$

$$+ ac - \frac{b^2}{a} > 0$$

$$\text{since } \left(x + \frac{b}{a}y\right)^2 > 0 \quad + \quad y^2 > 0$$

\Rightarrow

$$a > 0$$

$$ac - b^2 > 0.$$

$$7.3.1 \quad \dot{x} = x - y - x(x^2 + 5y^2)$$

$$\dot{y} = x + y - y(x^2 + y^2)$$

$$a) \quad A = \begin{bmatrix} 1 - 3x^2 - 5y^2 & -1 - 10xy \\ 1 - 2xy & 1 - x^2 - 3y^2 \end{bmatrix}$$

$$|A_{(0,0)} - \lambda I| = 0 \Rightarrow \begin{vmatrix} 1 - \lambda & -1 \\ 1 & 1 - \lambda \end{vmatrix} = 0$$

$$\Rightarrow (1 - \lambda)^2 + 1 = 0 \Rightarrow \lambda^2 - 2\lambda + 2 = 0$$

$$\Rightarrow \lambda = \frac{2 \pm \sqrt{4 - 4(2)}}{2} = 1 \pm i$$

$\therefore (0, 0)$ is an unstable spiral.

$$b) \quad r\dot{r} = x\dot{x} + y\dot{y}$$

$$= x^2 - \cancel{xy} - x^4 - 5x^2y^2 + \cancel{xy} + y^2 - x^2y^2 - y^4$$

$$= x^2 + y^2 - (x^4 + y^4 + 2x^2y^2) - 4x^2y^2$$

$$= r^2 - r^4 - 4r^2 \cos^2 \theta \cdot r^2 \sin^2 \theta$$

$$\Rightarrow \dot{r} = r - r^3 - r^3 \sin^2 2\theta$$

$$r^2 \dot{\theta} = x\dot{y} - y\dot{x}$$

$$= x^2 + \cancel{xy} - x^3y - x^2y^2 - \cancel{xy} + y^2 + x^3y + 5xy^3$$

$$= x^2 + y^2 + 4xy^3$$

$$= r^2 + 4r \cos \theta r^3 \sin^3 \theta$$

$$\Rightarrow \dot{\theta} = 1 + 4r^2 \cos \theta \sin^3 \theta$$

$$7.3.1 \quad c) \quad \dot{r} > 0 \Rightarrow r - r^3(1 + \sin^2 2\theta) > 0$$

$$\Rightarrow 1 > r^2(1 + \sin^2 2\theta)$$

$$\Rightarrow \frac{1}{\sqrt{1 + \sin^2 2\theta}} > r > 0$$

\therefore all trajectories have a radially outward component on the circle $r_1 = 1/\sqrt{2} - \varepsilon$

$$d) \quad \dot{r} < 0 \Rightarrow r - r^3(1 + \sin^2 2\theta) < 0$$

$$\Rightarrow 1 < r^2(1 + \sin^2 2\theta)$$

$$\Rightarrow \frac{1}{\sqrt{1 + \sin^2 2\theta}} < r$$

\therefore all trajectories have a radially inward component on the circle $r_2 = 1 + \varepsilon$

e) Since the region $1/\sqrt{2} - \varepsilon \leq r \leq 1 + \varepsilon$ is a trapping region that contains no fixed points, the Poincaré-Bendixon theorem implies there is a periodic orbit inside the region.

$$* \quad \dot{r} = 0 \Rightarrow r^2 = \frac{1}{1 + \sin^2 2\theta} \Rightarrow \dot{\theta} = 1 + 4 \frac{\cos \theta \sin^3 \theta}{1 + \sin^2 2\theta} = 0$$

$$\Rightarrow 1 + \sin^2 2\theta + 2 \sin 2\theta \sin^2 \theta = 0$$

$$\Rightarrow 1 + (\sin 2\theta + \sin^2 \theta)^2 - \sin^4 \theta = 0$$

$$\Rightarrow \sin^2 \theta (2 \cos \theta + \sin \theta)^2 + \cos^2 \theta (1 - \sin^2 \theta) = 0$$

$$\Rightarrow \sin \theta = \cos \theta = 0 \quad * \quad \therefore \dot{\theta} \neq 0$$

So no fixed points in the trapping region

7.3.4 $\dot{x} = x(1 - 4x^2 - y^2) - \frac{1}{2}y(1+x)$

$\dot{y} = y(1 - 4x^2 - y^2) + 2x(1+x)$

a) $A = \begin{bmatrix} 1 - 12x^2 - y^2 - \frac{1}{2}y & -2xy - \frac{1}{2} \\ -8xy + 2 + 4x & 1 - 4x^2 - 3y^2 \end{bmatrix}$

at $(0,0)$, $|A - \lambda I| = 0 \Rightarrow \begin{vmatrix} 1-\lambda & -1/2 \\ 2 & 1-\lambda \end{vmatrix} = 0$

$\Rightarrow (1-\lambda)^2 + 1 = 0 \Rightarrow \lambda^2 - 2\lambda + 2 = 0$

$\Rightarrow \lambda = \frac{2 \pm \sqrt{4 - 4(2)}}{2} = 1 \pm i$

$\Rightarrow (0,0)$ is an U.S. spiral

b) $V = (1 - 4x^2 - y^2)^2$ use \dot{V} to show trajectories approach $4x^2 + y^2 = 1$.

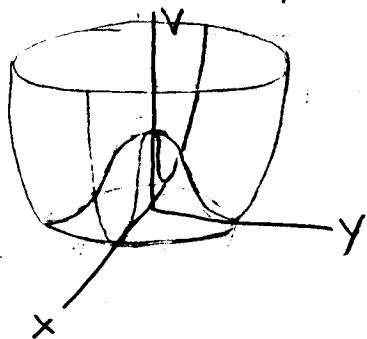
$\dot{V} = 2(1 - 4x^2 - y^2)(-8xx' - 2yy')$

$= 2(1 - 4x^2 - y^2)(-8x^2 + 32x^4 + 8x^2y^2 + 4xy + 4x^2y - 2y^2 + 8x^2y^2 + 2y^4 - 4xy - 4x^2y)$

$= -4(1 - 4x^2 - y^2)(4x^2 + y^2 - 8x^2y^2 - 16x^4 - y^4)$

$= -4(1 - 4x^2 - y^2)^2(4x^2 + y^2) < 0$ unless $x=y=0$ or $4x^2 + y^2 = 1$

• V decreases along trajectories until it reaches the ellipse $4x^2 + y^2 = 1$ (b/c $(0,0)$ is a max. of V).



So all trajectories approach the ellipse -- except $(0,0)$.

$$? 5.4 \quad \ddot{x} + u f(x) \dot{x} + x = 0$$

$$f(x) = \begin{cases} -1 & |x| < 1 \\ 1 & |x| \geq 1 \end{cases}$$

$$a) \quad \int f(x) dx = \begin{cases} \int -dx & |x| \leq 1 \\ \int dx & x \geq 1 \\ \int dx & x \leq -1 \end{cases}$$

$$F(x) = \int f(x) dx$$

$$F(x) = \begin{cases} -x & |x| \leq 1 \\ x+c & x \geq 1 \\ x+c & x \leq -1 \end{cases}$$

to get c want

$$x+c = -x \quad \text{at } x=1 \\ \Rightarrow c = -2$$

$$x+c = -x \quad \text{at } x=-1 \\ \Rightarrow c = 2$$

S₀

$$F(x) = \begin{cases} x+2, & x \leq -1 \\ -x, & |x| \leq 1 \\ x-2, & x \geq 1 \end{cases}$$

$$\ddot{x} + \mu \dot{x} f(x) = \frac{d}{dt} (\dot{x} + \mu F(x))$$

$$w = \dot{x} + \mu F(x)$$

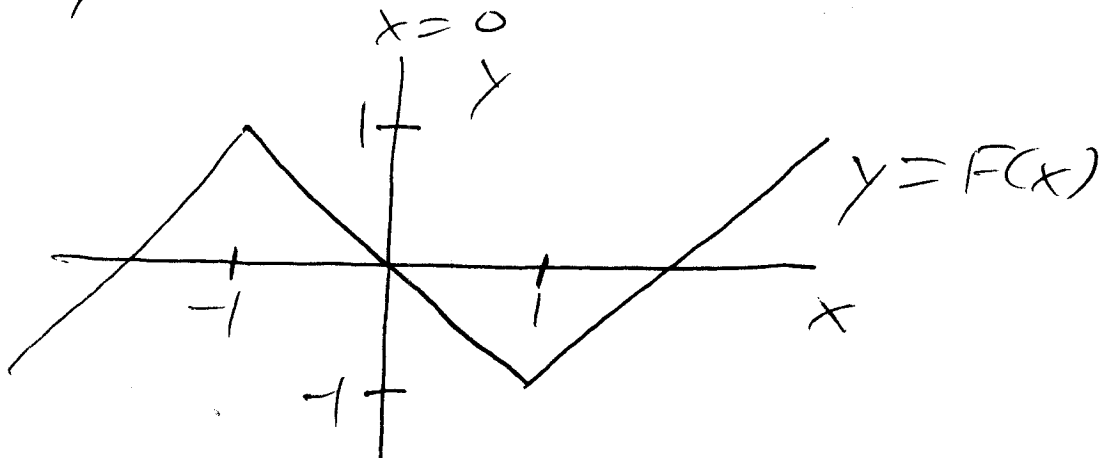
$$\Rightarrow \dot{w} = -x$$

$$\begin{cases} \dot{x} = w - \mu F(x) \\ \dot{w} = -x \end{cases}$$

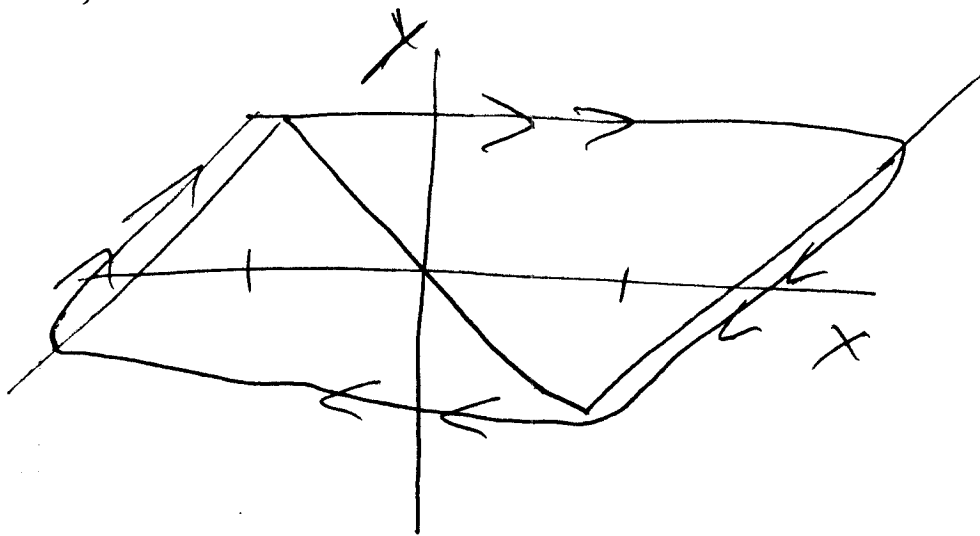
$$\text{let } y = w/\mu$$

$$\Rightarrow \begin{cases} \dot{x} = \mu (y - F(x)) \\ \dot{y} = -x/\mu \end{cases}$$

b)



c) The oscillation would look like



We can justify this picture by saying:

suppose that the initial condition is not too close to the piece-wise linear nullcline, i.e.

suppose $y - F(x) \sim O(1)$. Then

$$\begin{cases} \dot{x} = \mu(y - F(x)) \\ \dot{y} = -x/\mu \end{cases}$$

implies that $|\dot{x}| \sim O(\mu) \gg 1$

and $|\dot{y}| \sim O(\mu^{-1}) \ll 1$

hence, the velocity is enormous in the horizontal direction and tiny in the vertical direction, so trajectories move practically horizontally.

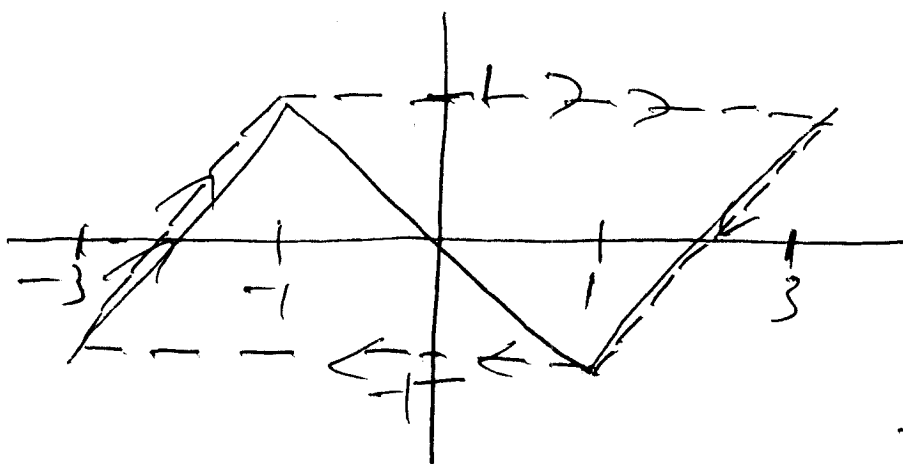
If the initial condition is above the nullcline, then $y - F(x) > 0$ and $\dot{x} > 0$; thus the trajectory moves sideways toward the nullcline. If the initial condition is below the nullcline, then $y - F(x) < 0$ and $\dot{x} < 0$. However, once the trajectory gets so close that $y - F(x) \sim O(\epsilon^{-2})$, then \dot{x} and \dot{y} become comparable and the trajectory crosses the nullcline vertically.

d) The period of the oscillations is the integral of the two linear functions $x - 1$ and $x + 1$.
~~to find the limits of integration~~

This is true since

$y \approx F(x)$ on the slow branches of the nullcline

To find the limits of integration



Find x when

$$x - 1 = 1$$

$$\Rightarrow x = 2$$

$$x + 1 = -1$$

$$\Rightarrow x = -2$$

next,

$$y \approx F(x)$$

$$\dot{y} \approx F'(x) \frac{dx}{dt} = (1) \frac{dx}{dt}$$

$$\dot{y} = -x/m$$

$$\Rightarrow -x/m = \frac{dx}{dt}$$

$$\Rightarrow dt \approx -\frac{m}{x} dx$$

$$T \approx 2 \int_3^1 -\frac{m}{x} dx = 2 \int_1^3 \frac{m}{x} dx$$

$$= 2m [\ln(x)]_1^3$$

$$= 2m \ln(3)$$